

KBM Nonlinear Dynamics and First-Principles-Based Features in Deep Learning Algorithm for Predicting Disruptions

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Outline

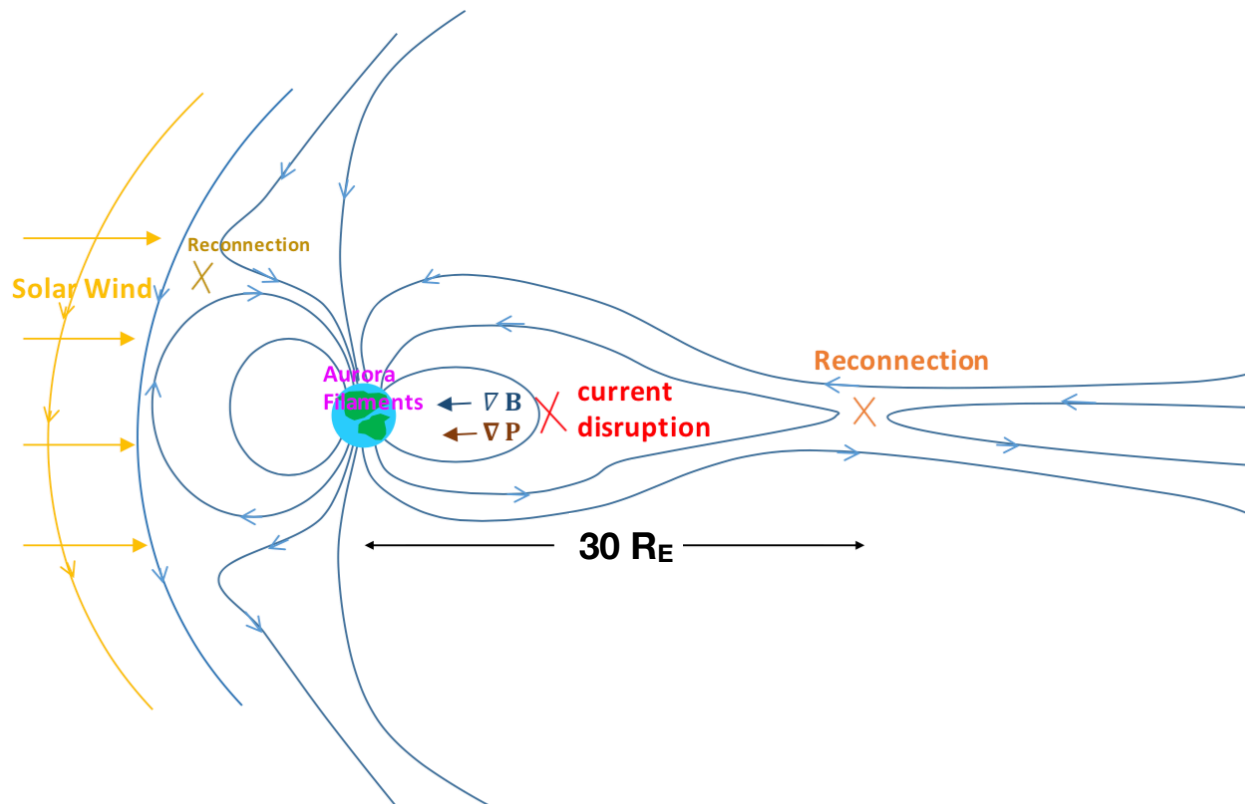
Nonlinear KBM

- Introduction
- KBM saturation mechanism in Cyclone Base Case
- KBM nonlinear dynamics in DIII-D pedestal

Future work

- Motivation
- Feeding First-Principles-Based code output to AI

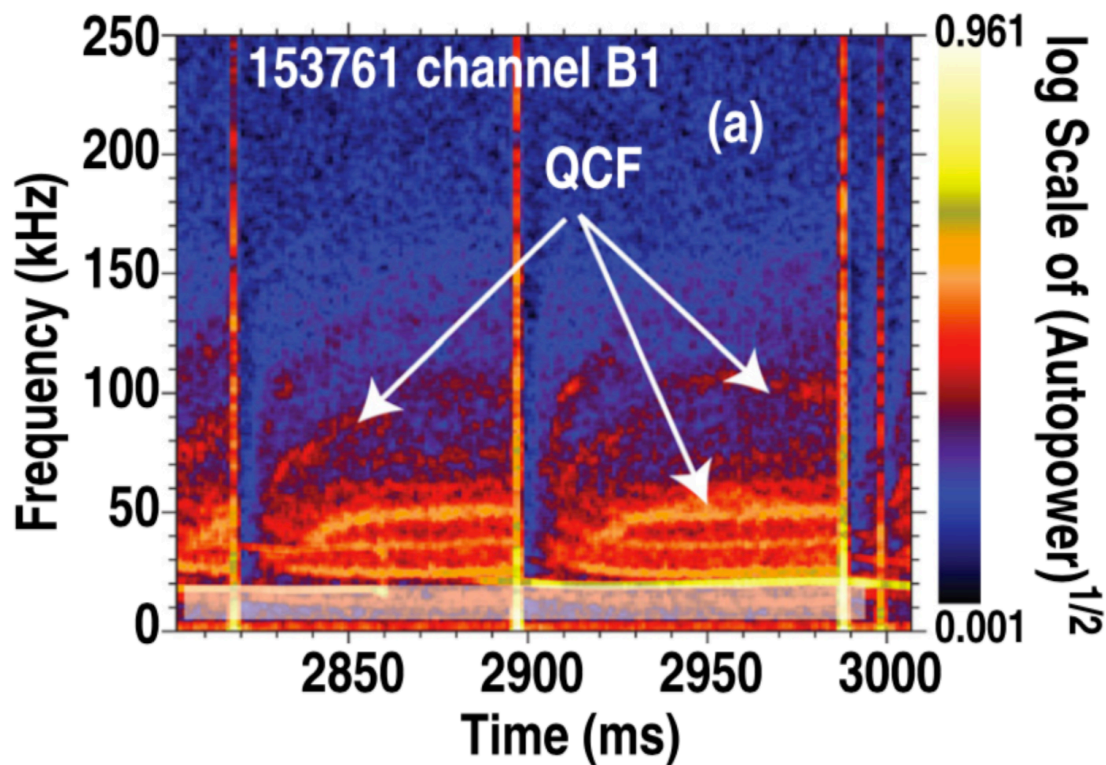
Ballooning mode as a candidate for explosive behavior in substorm onset



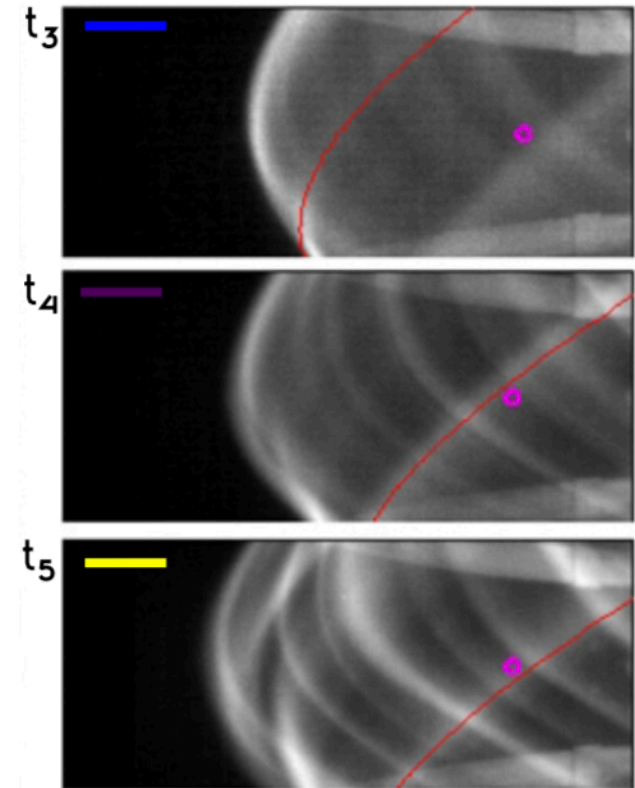
A Sketch of substorm onset.

Ballooning modes are proposed as a candidate for triggering explosive behavior in substorm onset. [Bhattacharjee et al., 1998; Gohil et al., 1988; Raeder et al. 2010]

High frequency coherent fluctuations are related to ELM activities in various tokamak devices

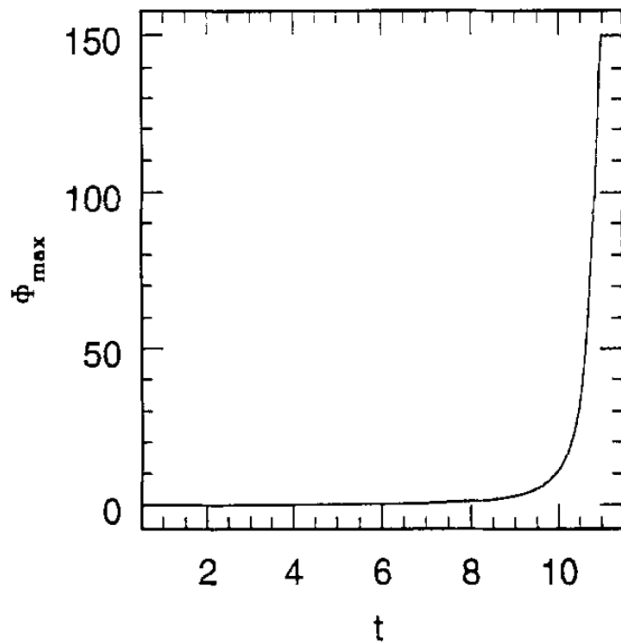


Quasi-coherent fluctuations and ELM crashes in DIII-D. Figure reproduced from A. Diallo, *Phys. Plasmas* 22, 056111 (2015). Inter-ELM magnetic fluctuations spectrograms as measured using the Mirnov coils.

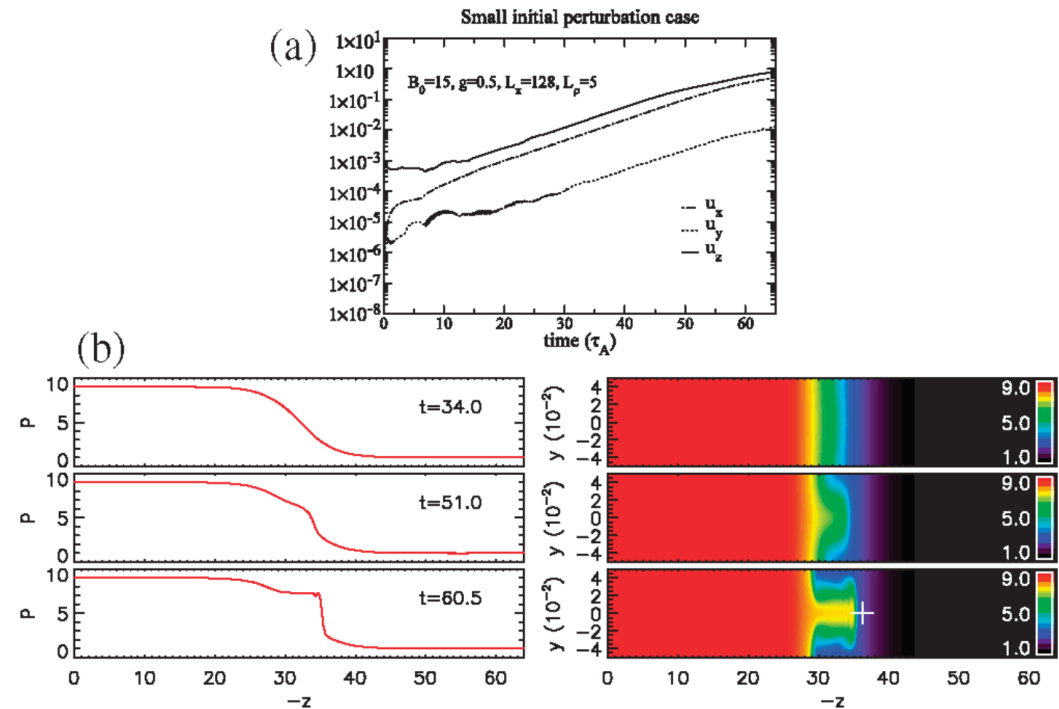


Filamentary structures of ELM on MAST. Figure reproduced from R. Scannell, *PPCF*, 49, 1431, (2007). Filaments observed during an ELM by the Photron camera. The camera frames are separated by 16 μ s

Explosive growth of IBM and the intermediate regime in MHD simulations



Finite time singularity of the flow poloidal gradient in ideal ballooning mode near marginal instability.
Figure reproduced from *S. C. Cowley, Physics Reports 283 (1997)*



Exponential nonlinear growth of weakly unstable ballooning modes in full MHD simulations.

Intermediate regime is formulated to describe the nonlinear behavior.

Figure reproduced from *P. Zhu, PRL 96, 065001 (2006)*

Gyrokinetic simulation model

Gyrokinetic ion model.

Solve Vlasov equation directly.

Conservative scheme for electrons.

Separate adiabatic $\delta f_e^{(a)}$ and non-adiabatic part δh_e .

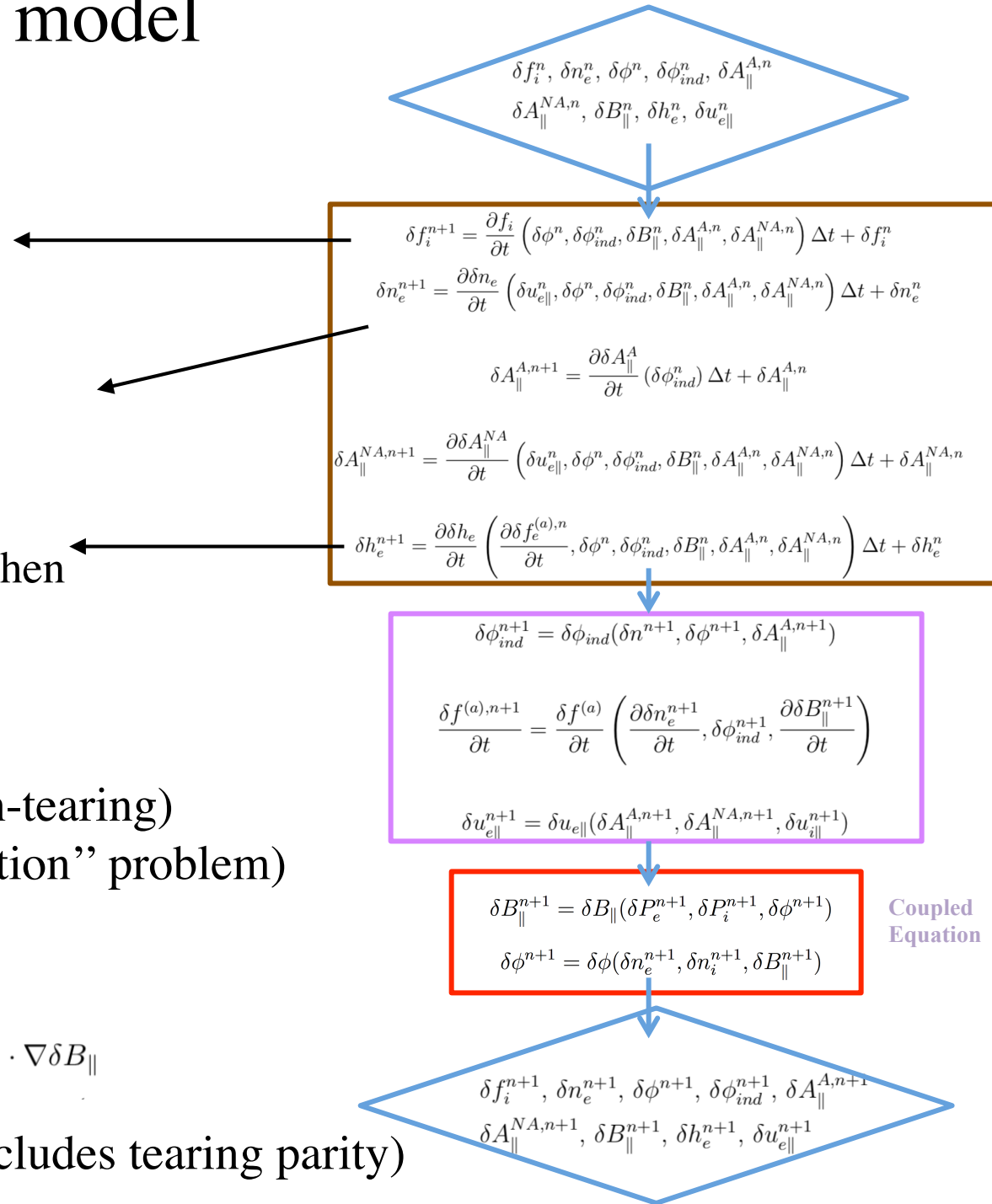
Solve δn_e from continuity equation, then $\delta f_e^{(a)}$ analytically.

$$\frac{\partial \delta A_{\parallel}^A}{\partial t} = c b_0 \cdot \nabla \phi_{ind}. \quad \text{Adiabatic (non-tearing)} \\ \text{(No ``cancellation'' problem)}$$

$$\frac{1}{c} \frac{\partial \delta A_{\parallel}^{na}}{\partial t} = \frac{\delta \mathbf{B}_{\perp}}{B_0} \cdot \nabla \delta \phi_{ind} \\ - \frac{m_e}{n_0 e^2} \nabla \cdot \left(\delta u_{\parallel e} \frac{c P_{e0} \mathbf{B}_0 \times \nabla \delta B_{\parallel}}{B_0^3} \right) + \frac{P_{e0}}{e n_0} \frac{\delta \mathbf{B}_{\perp}}{B_0^2} \cdot \nabla \delta B_{\parallel}$$

Non-adiabatic (includes tearing parity)

GTC Flow Chart-Conservative



Gyrokinetic simulation model

Gyrokinetic Poisson Equation And Perpendicular Ampere's Law

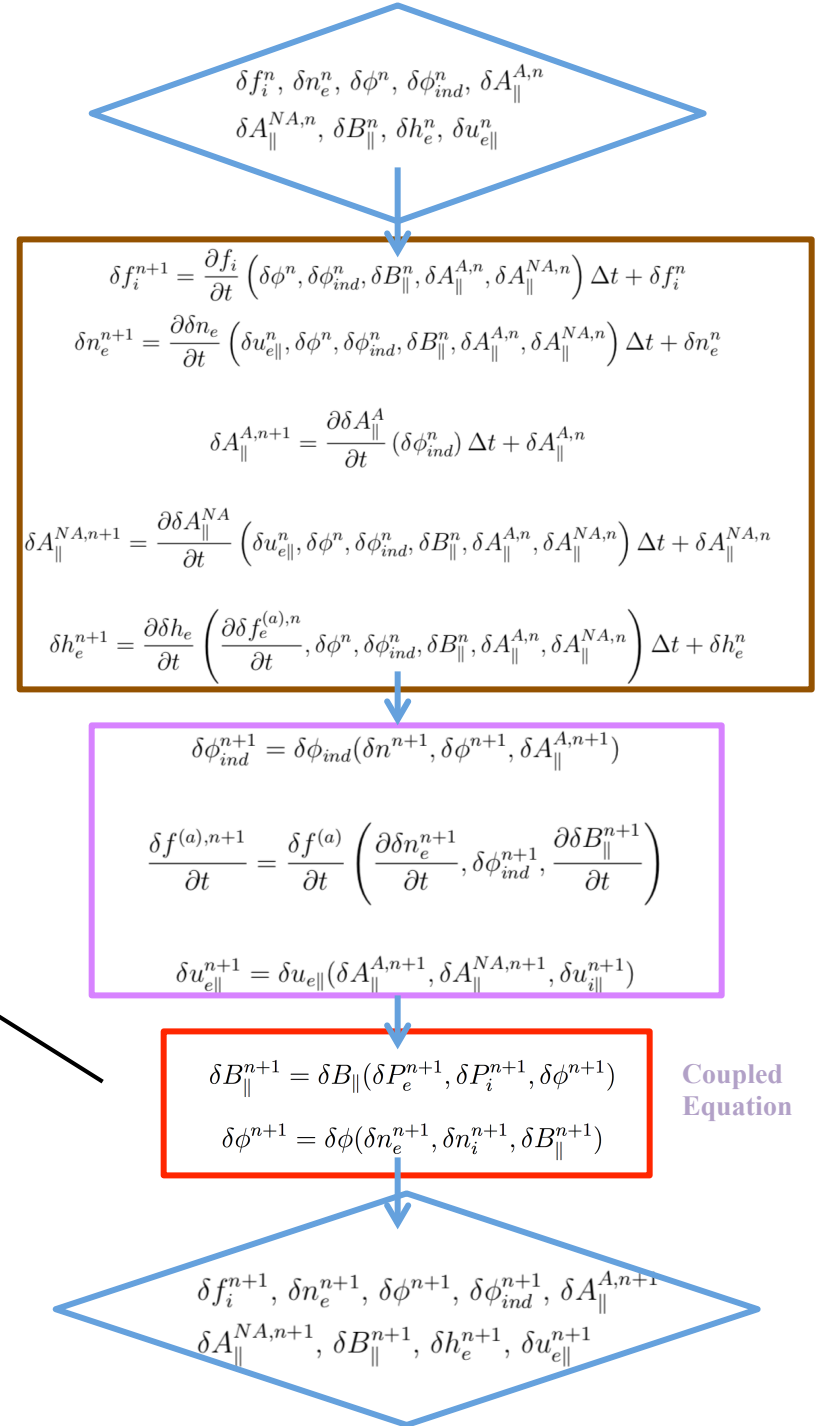
$$\frac{Z_i^2 n_i}{T_i} (\delta\phi - \tilde{\delta\phi}) - \frac{1}{B_0} (Z_i n_{i0} \{\delta B_{\parallel}\}_i - e n_{e0} \{\delta B_{\parallel}\}_e) = Z_i n_i - e n_e,$$

$$\begin{aligned} & \frac{\delta B_{\parallel} B_0}{4\pi} + 2\pi\Omega_e^2 \int d\mu dv_{\parallel} \left[B_0 \left\langle \int_0^{\rho_e} F_e^{\text{gyro}} r dr \right\rangle \right. \\ & \quad \left. + \frac{f_M}{\rho_e^2} \left\langle \int_0^{\rho_e} \left\langle \int_0^{\rho_e} \delta B_{\parallel} r' dr' \right\rangle r dr \right\rangle \right] \\ & = -2\pi\Omega_i^2 \int d\mu dv_{\parallel} \left[B_0 \left\langle \int_0^{\rho_i} \left(F_i^{\text{gyro}} + \frac{e \langle \delta\phi \rangle - e \delta\phi}{T_i} F_M \right) r dr \right\rangle \right. \\ & \quad \left. + \frac{F_M}{\rho_i^2} \left\langle \int_0^{\rho_i} \left\langle \int_0^{\rho_i} \delta B_{\parallel} r' dr' \right\rangle r dr \right\rangle \right]. \end{aligned} \quad (12)$$

$$\tilde{\delta\phi} = \int d\mathbf{v} \int d\mathbf{R} \langle \delta\phi \rangle(\mathbf{R}) F_M(\mathbf{R}, v_{\parallel}, \mu, t) \delta(\mathbf{x} - \mathbf{R} - \rho),$$

$$\begin{aligned} \{\delta B_{\parallel}\}_s &= \frac{m\Omega_s^2}{2\pi T_s} \int d\mathbf{v} \int d\mathbf{R} \int_0^{\rho} r' dr' \int_0^{2\pi} d\zeta' \int d\mathbf{x}' \delta B_{\parallel}(\mathbf{x}') \\ &\quad \times \delta(\mathbf{x}' - \mathbf{R} - \rho) F_M \delta(\mathbf{x} - \mathbf{R} - \rho). \end{aligned}$$

GTC Flow Chart-Conservative



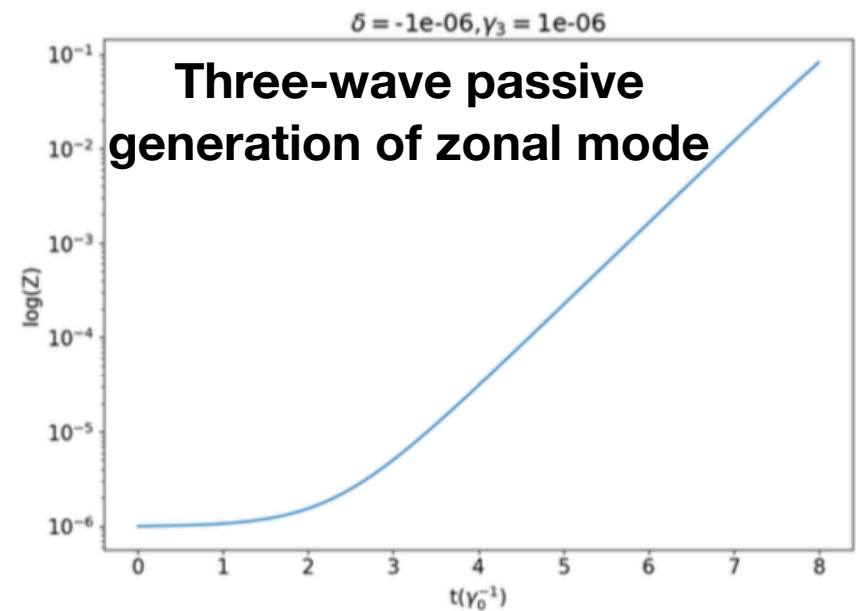
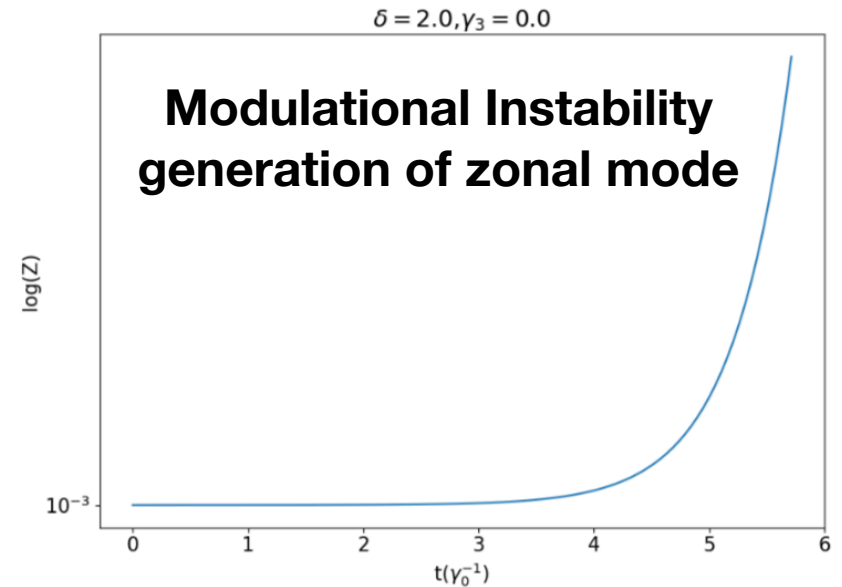
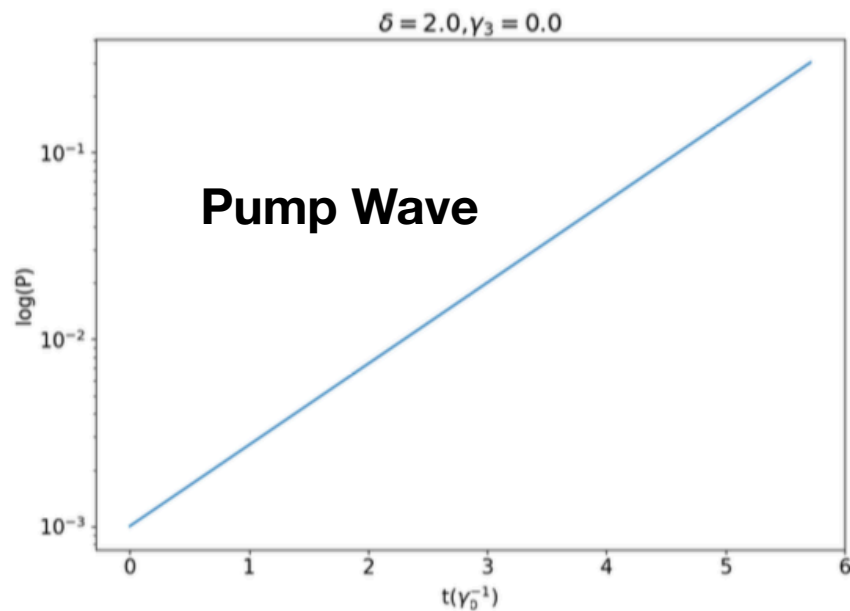
Spontaneous zonal flow generation

$$\phi_0(\mathbf{r}, t) = e^{-i(n\phi + \omega_0 t)} \sum_m \Phi_0(m - nq) e^{im\theta},$$

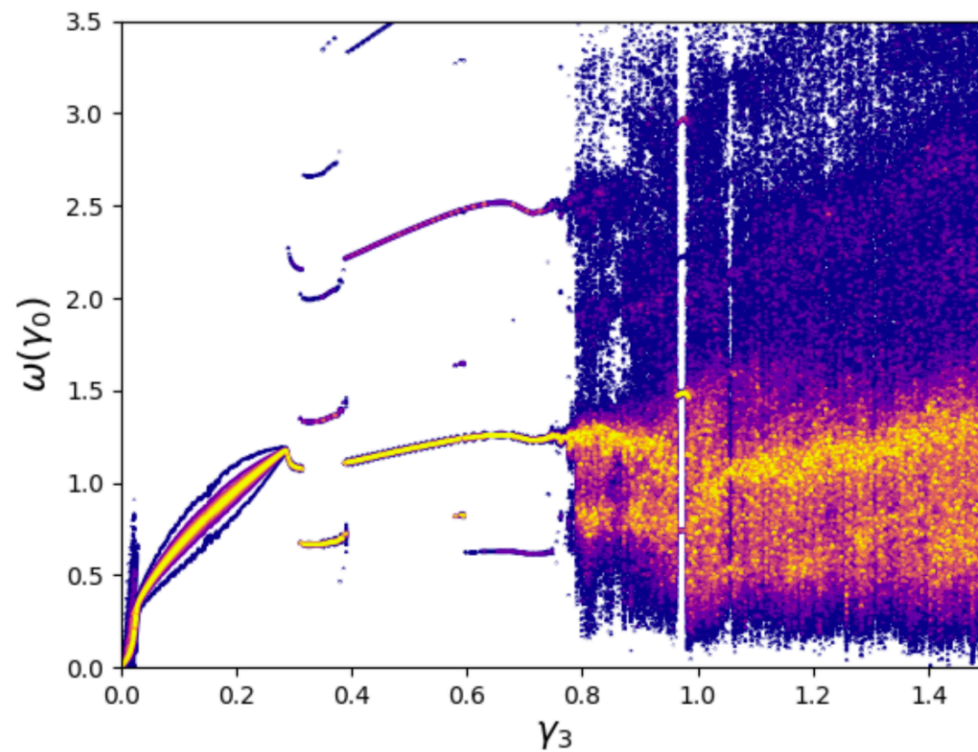
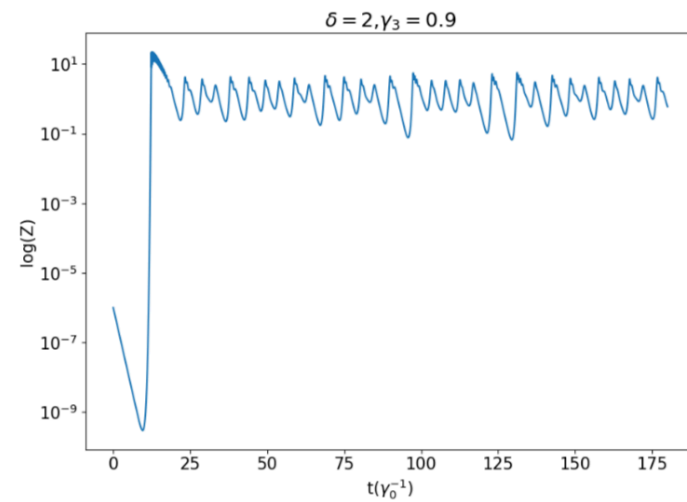
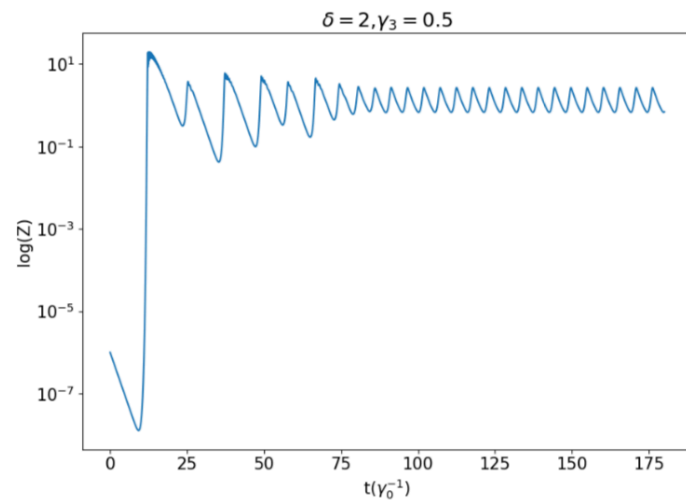
$$\delta\phi_{\pm} = e^{i(\mp n\phi - (\omega_z \pm \omega_0)t + K_z r)} \sum_m \Phi_{\pm}(m - nq) e^{im\theta},$$

$$\delta\phi_z = \Phi_z e^{i(K_z r - \omega_z t)} + \text{c.c.}$$

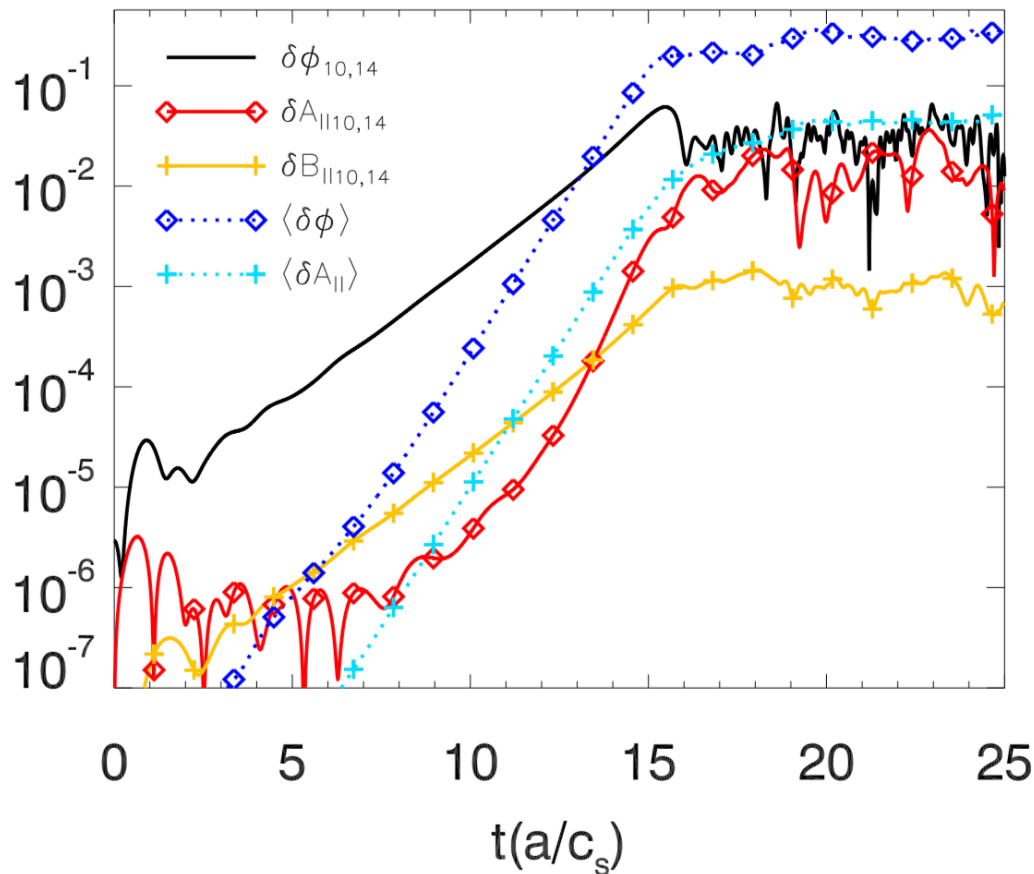
$$\text{Gamma}_3 \sim \nu_z = (1.5 \epsilon \tau_{ii})^{-1}$$



Route to Chaos



Zonal fields Saturation of KBM in Cyclone Base Case in global Gyrokinetic simulations



KBM nonlinear time history

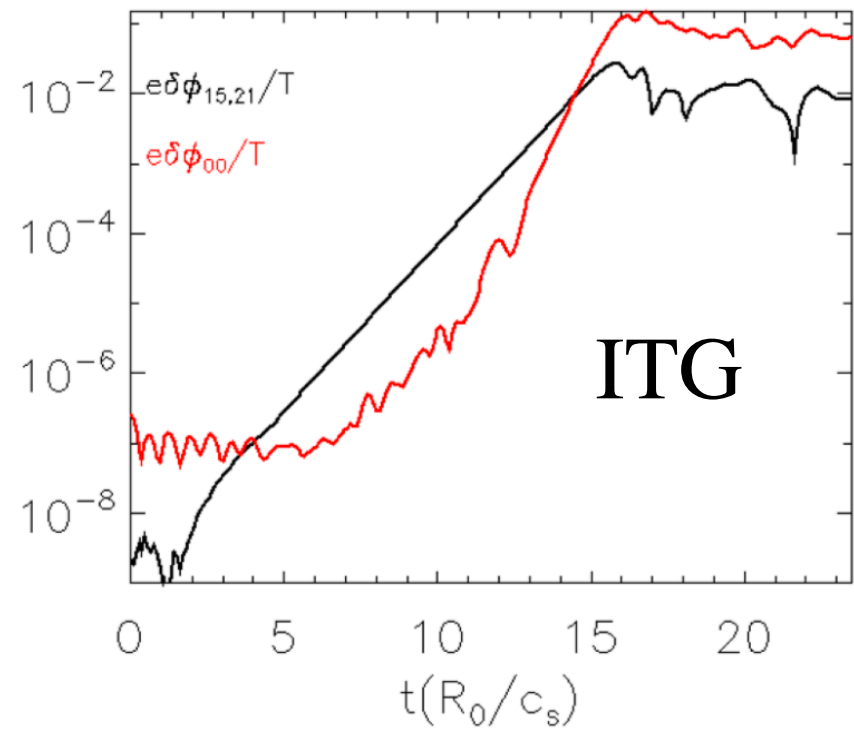
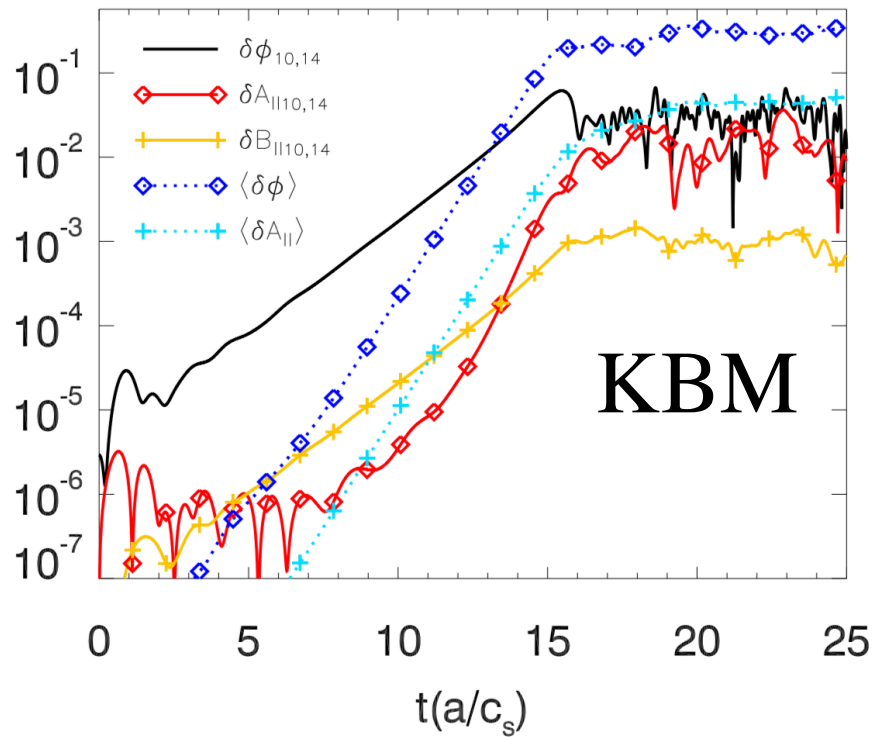
Fixed equilibrium ∇P_0 .

IBM unstable

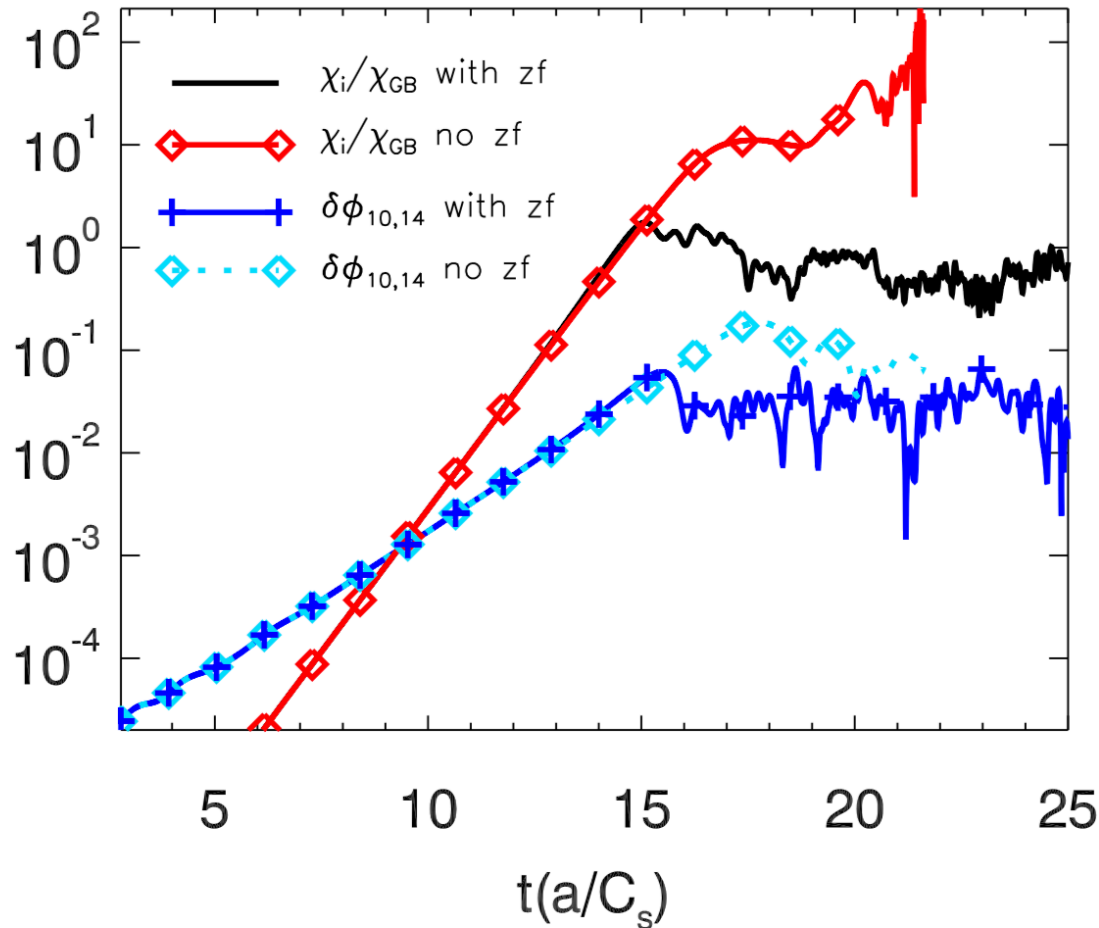
CBC parameter: $R_0 = 83.5\text{cm}$, $a/R_0 = 0.357$. At $r = 0.5a$, $B_0 = 2.01\text{T}$, $T_e = 2223\text{eV}$, $R_0/L_T = 6.9$, $R_0/L_n = 2.2$, $q = 1.4$, $\beta_e = 2\%$.

First order $s - \alpha$ model with circular cross-section $\theta = \theta_0 - \varepsilon \sin \theta_0 + O(\varepsilon^2)$

Comparison of ITG and KBM zonal flow generation



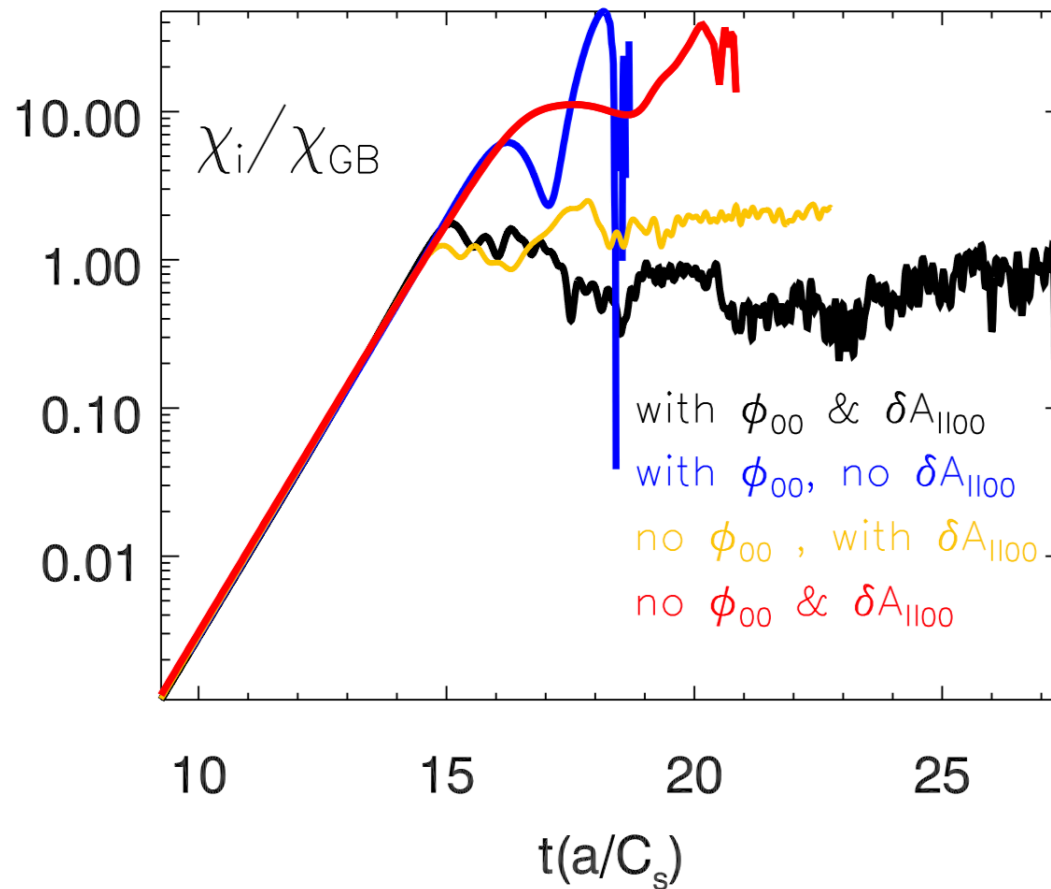
Zonal fields regulate mode saturation



$$\chi_i = \frac{1}{n_0 \nabla T_i} \int d\mathbf{v} \left(\frac{1}{2} m_i v^2 - \frac{3}{2} T_i \right) v_r \delta f$$

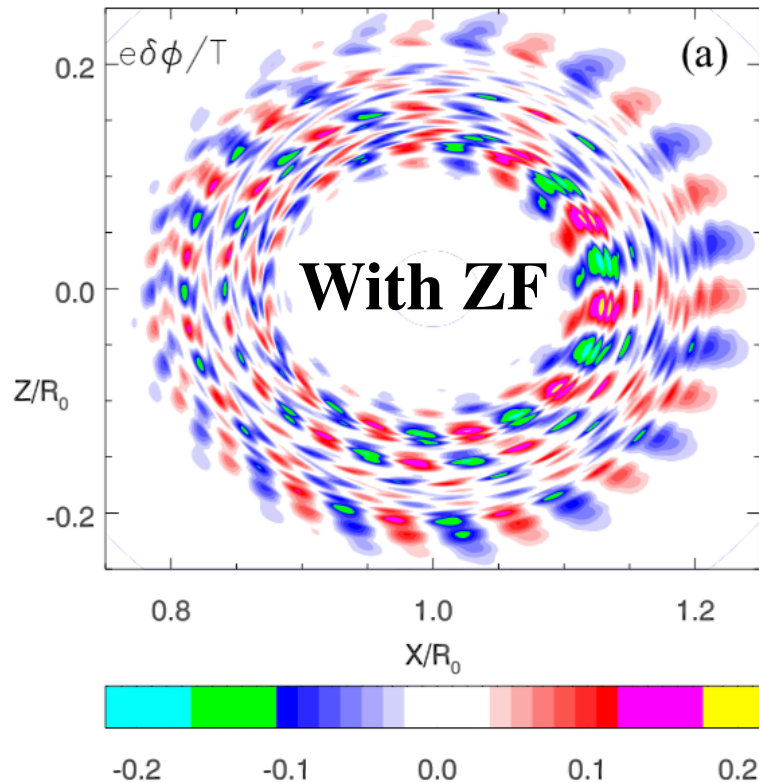
Comparison of simulation with and w/o zonal fields.
In the case w/o zonal fields, the mode saturates at almost one order of magnitude higher

Both the Zonal flow and Zonal Current are important

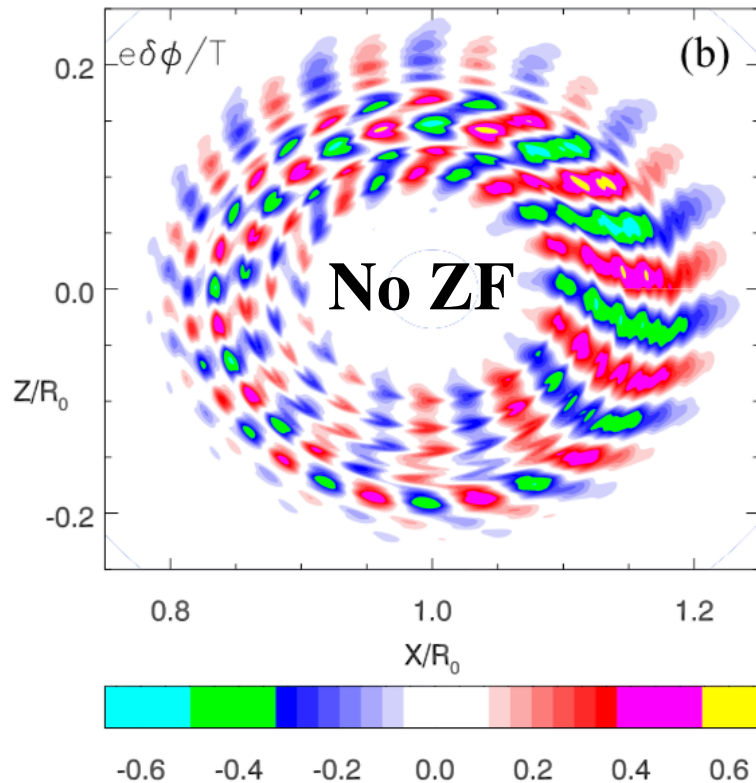


Comparison of 4 simulation with and w/o zonal current/ zonal flow respectively.

Zonal fields breaking of mode structure



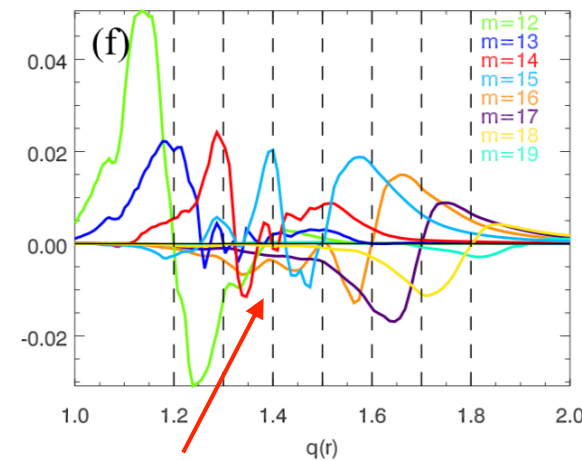
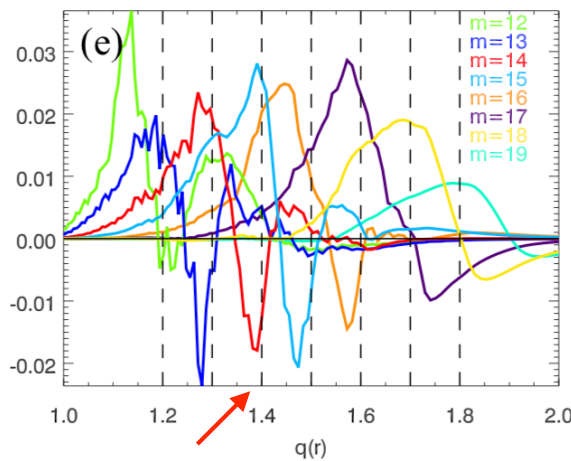
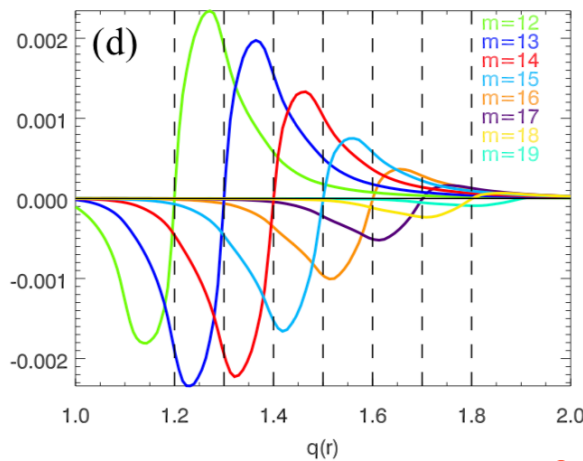
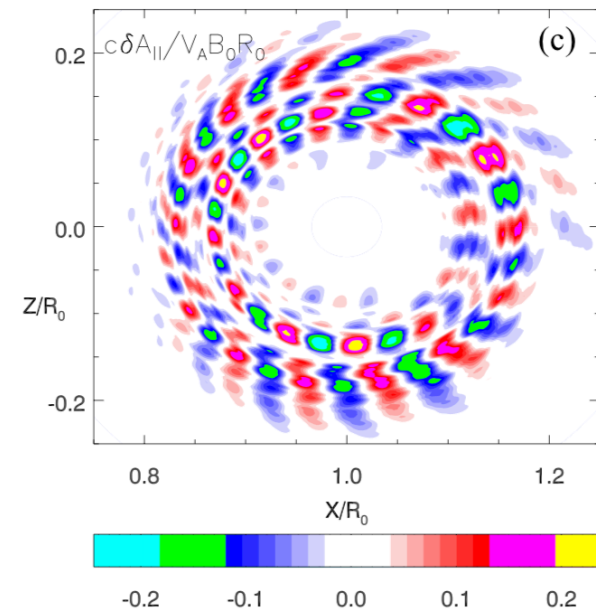
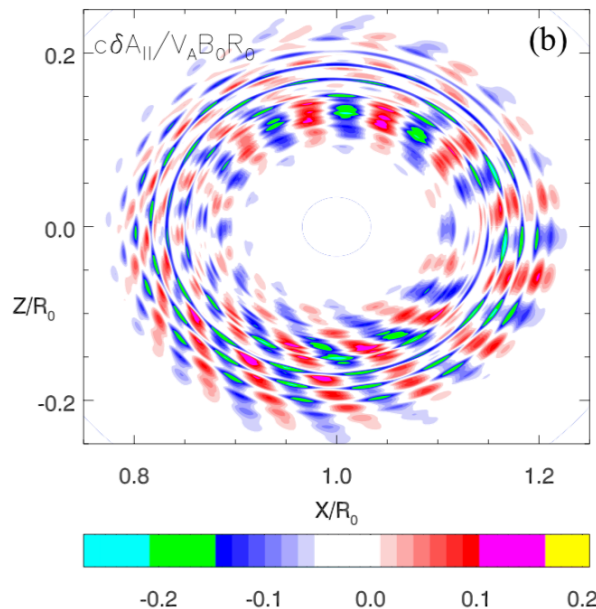
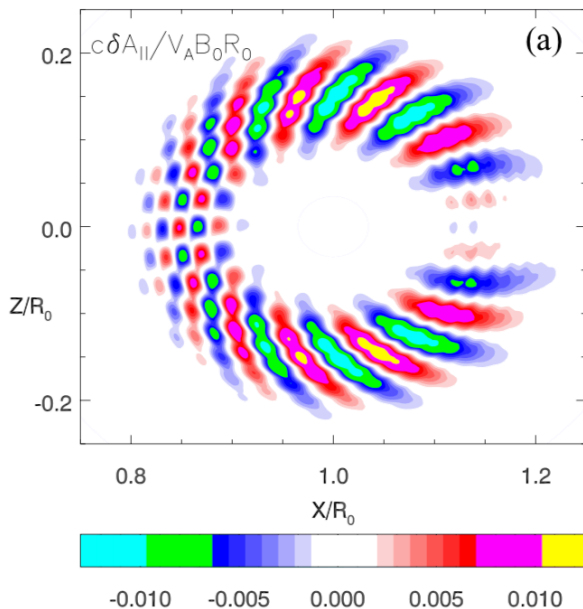
Broken radial filaments with self-consistently generated zonal flow and zonal current.



Macroscale radial filaments in the simulation without zonal fields.

Generation of high n localized current sheet

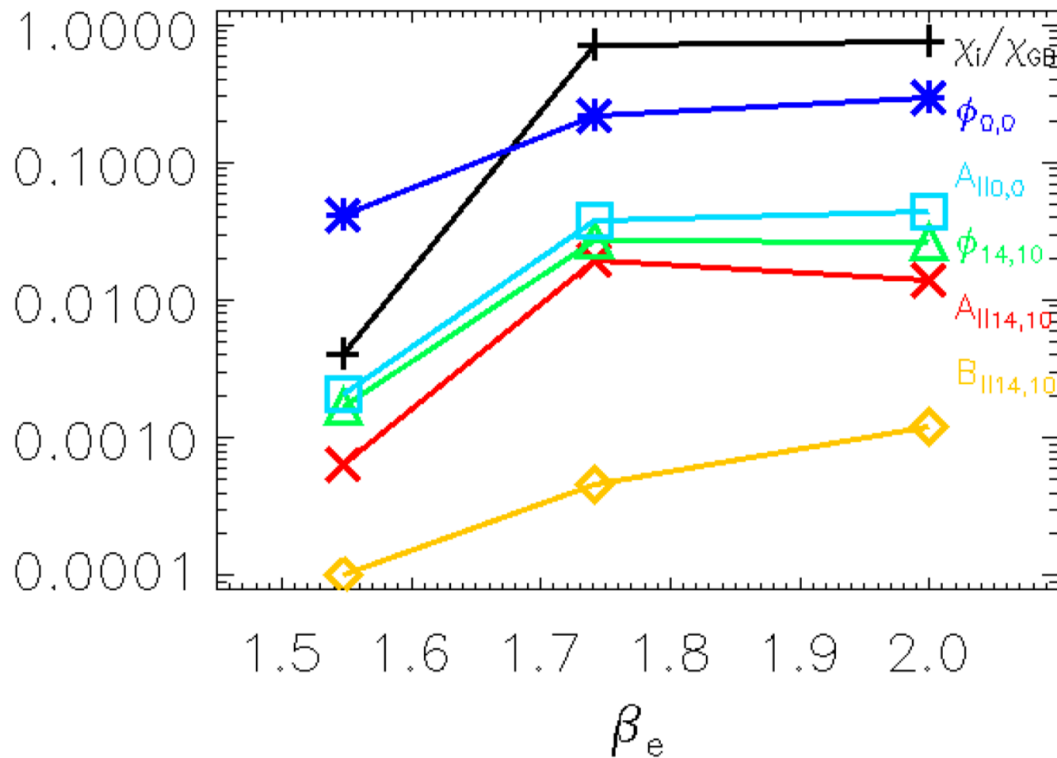
Non-tearing mode nonlinearly can induce tearing mode



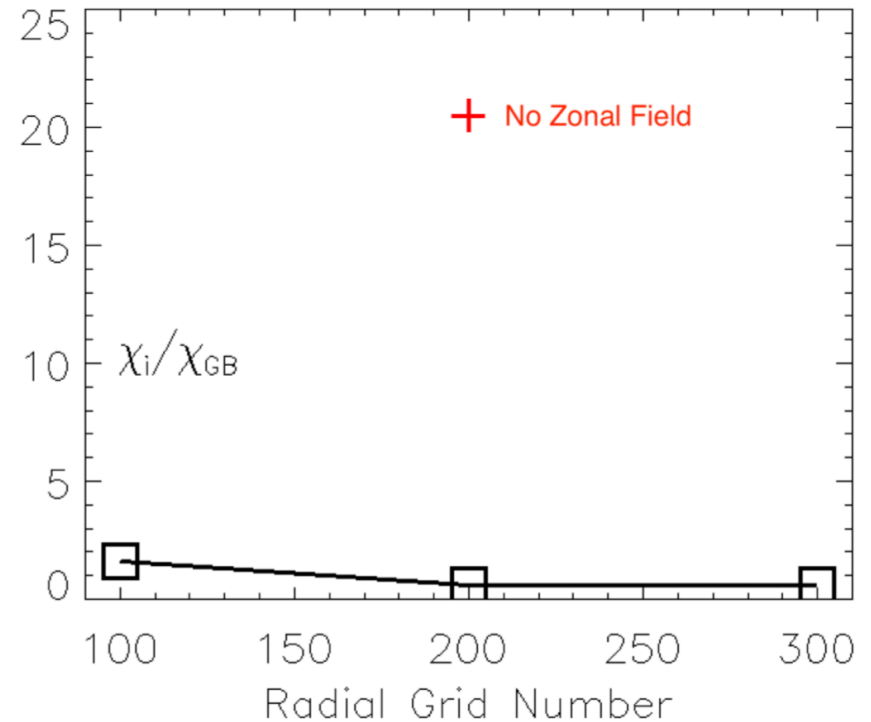
$\delta A_{||}$ finite at Rational Surface $\delta A_{||}=0$ at RS

Saturation mechanism not sensitive to β_e

Saturation amplitude



Similar saturation features for lower β_e

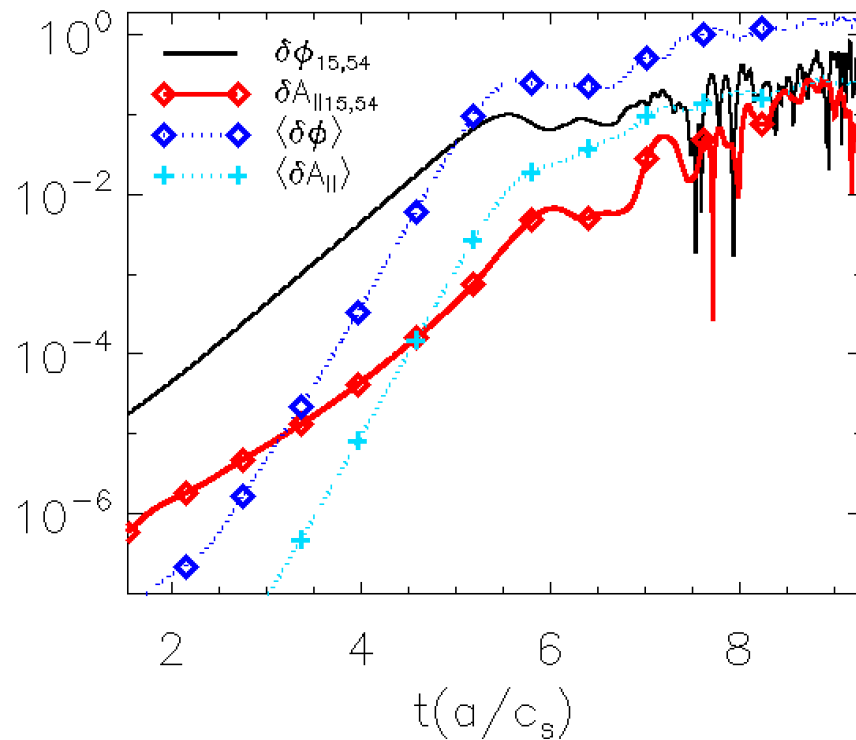


Convergence study

KBM in DIII-D pedestal steep gradient region

Table 4.1: Parameter for DIII-D shot #145701 at the steep gradient region at $\psi_n = 0.985$

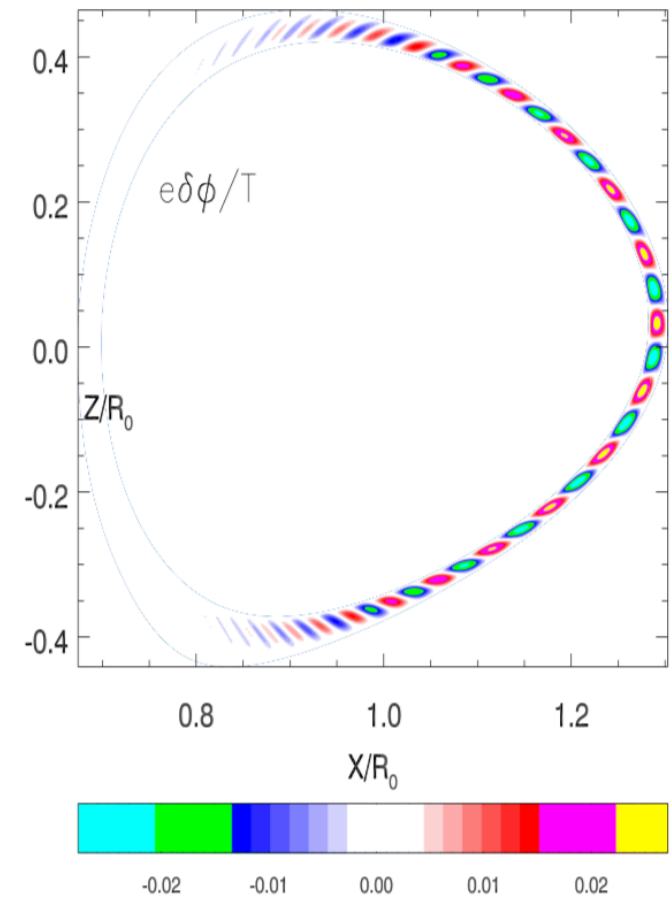
T_e	T_i	n_e	L_n/L_{Ti}	R_0/L_n	R_0/L_{Te}	R_0/L_{Ti}	q	β_e
197ev	396ev	$2.48 \times 10^{19} m^{-3}$	0.13	64	144	8	0.36	0.7%



Time history of nonlinear mode evolution for KBM in DIII-D edge

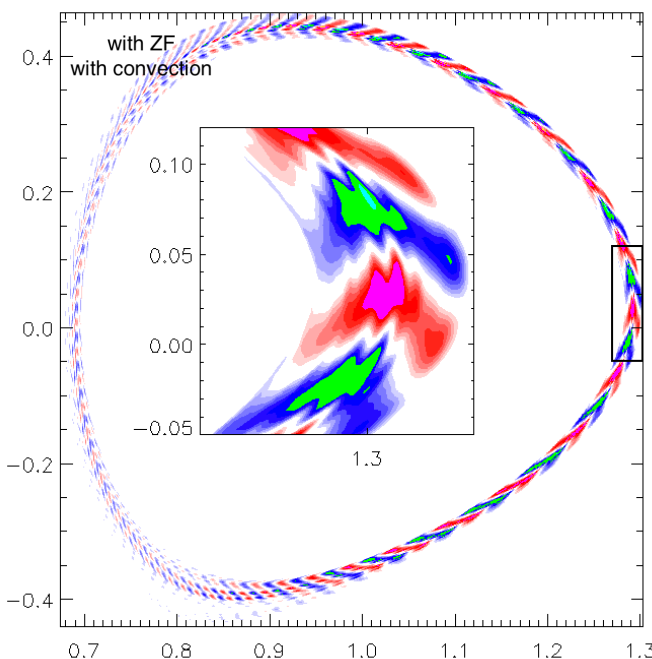
No equilibrium E_r shear

Collisionless

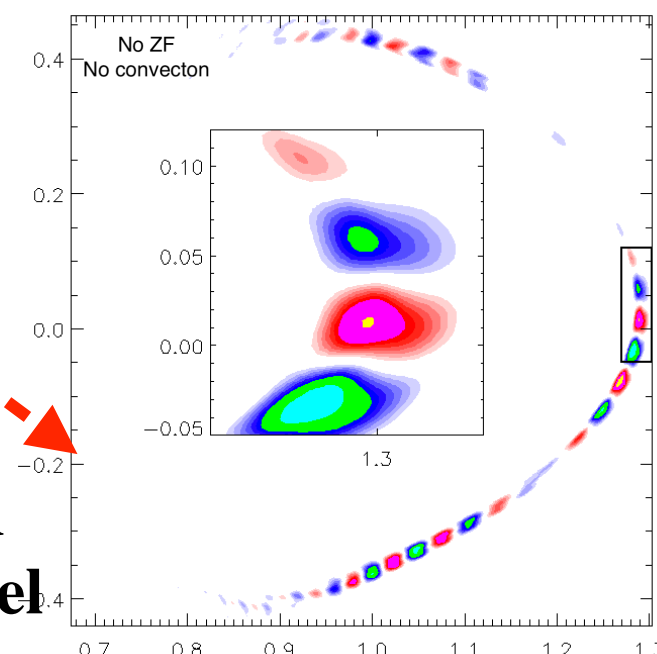
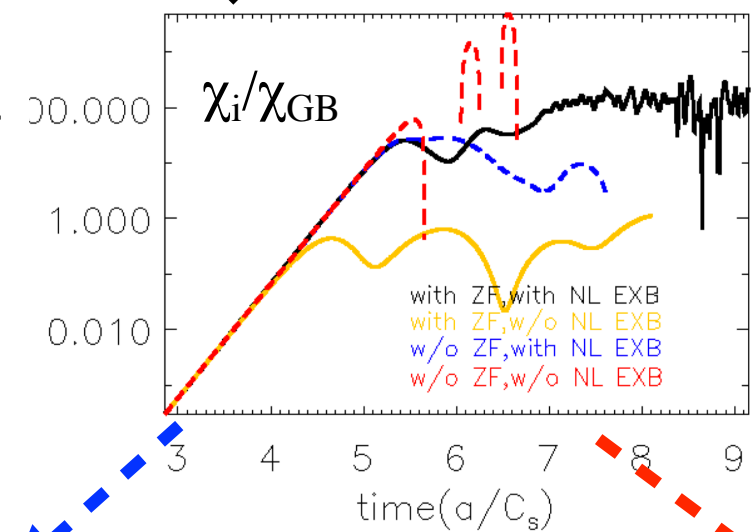
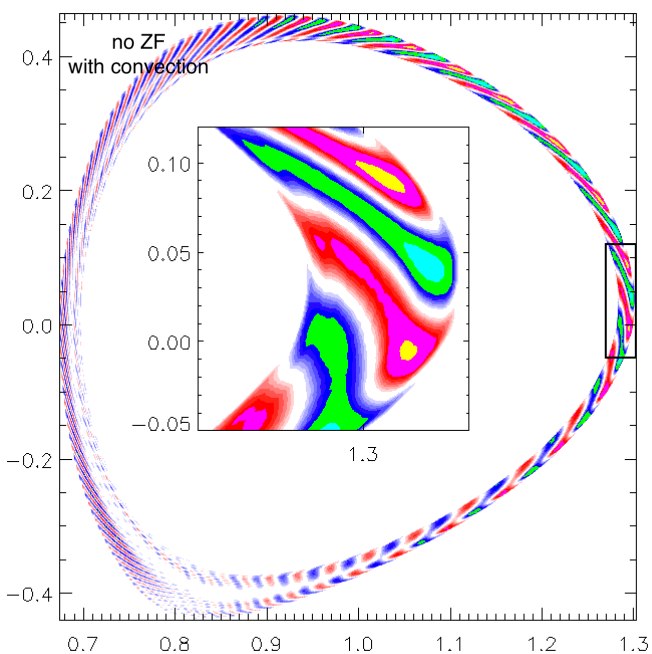
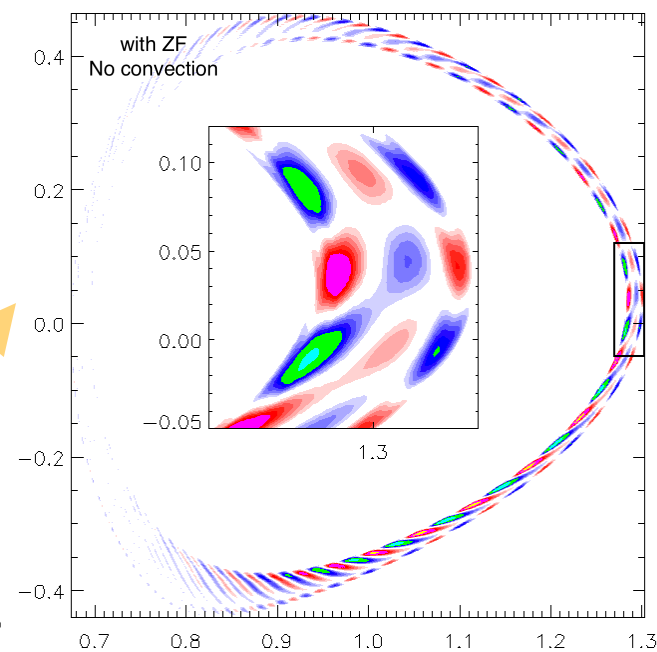


Linear mode structure

Competing effects of ZF and the ExB convection

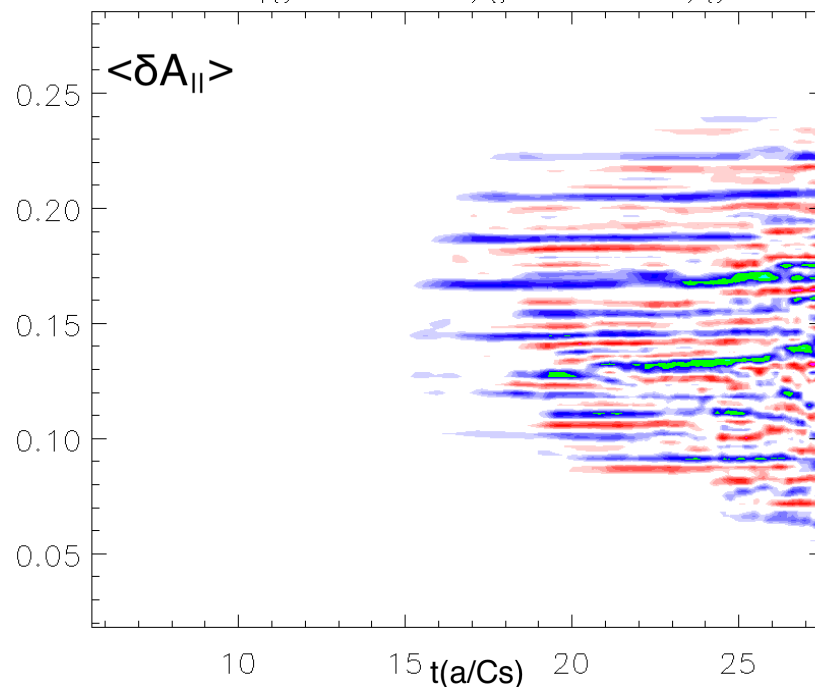
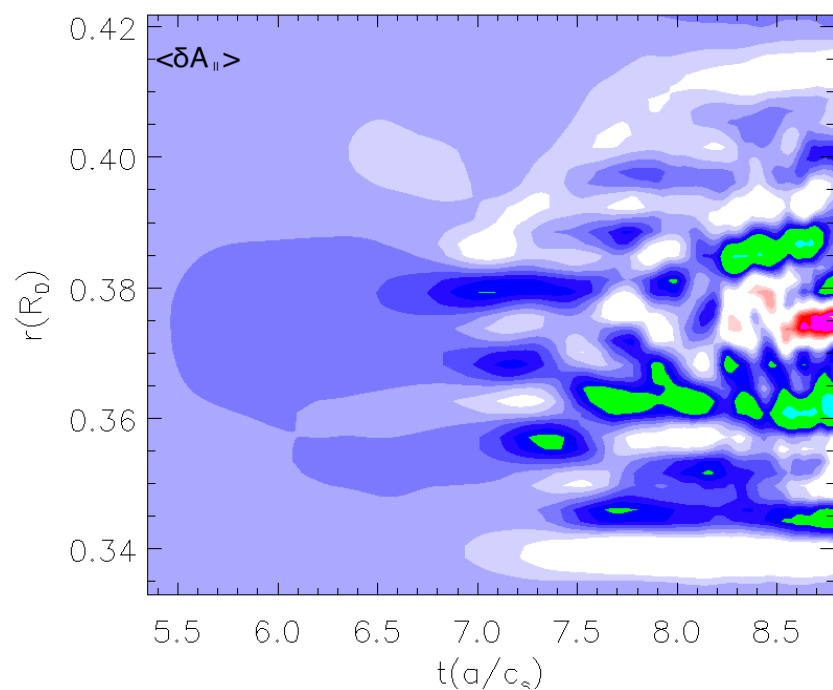
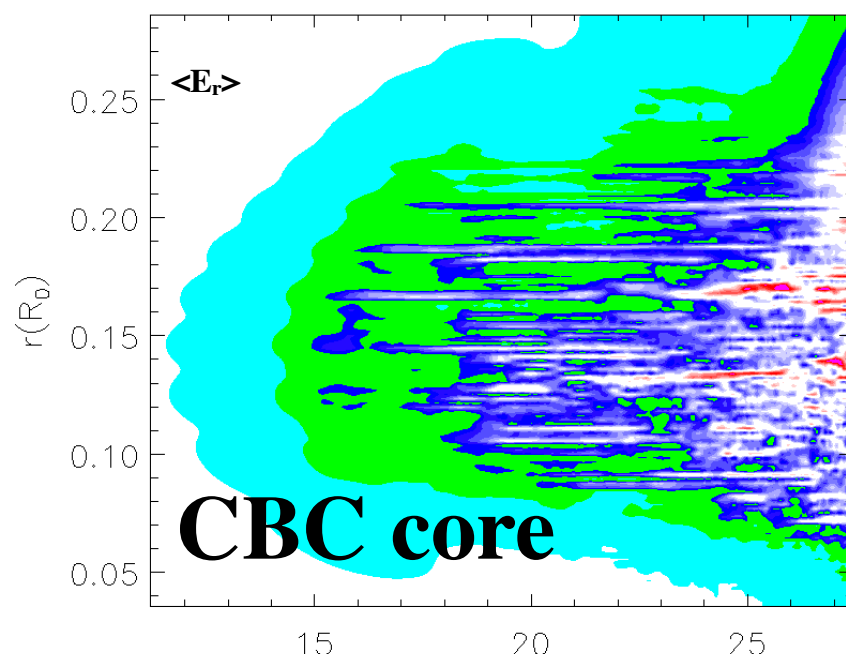
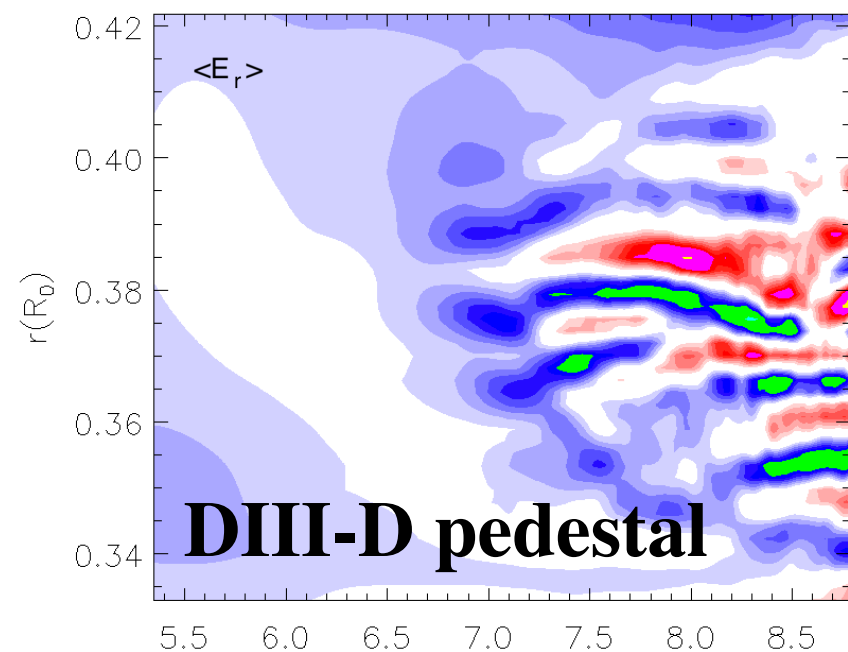


Nonlinear KBM $\delta\phi$
poloidal contour



Mode shearing of ZF and
convection partially cancel

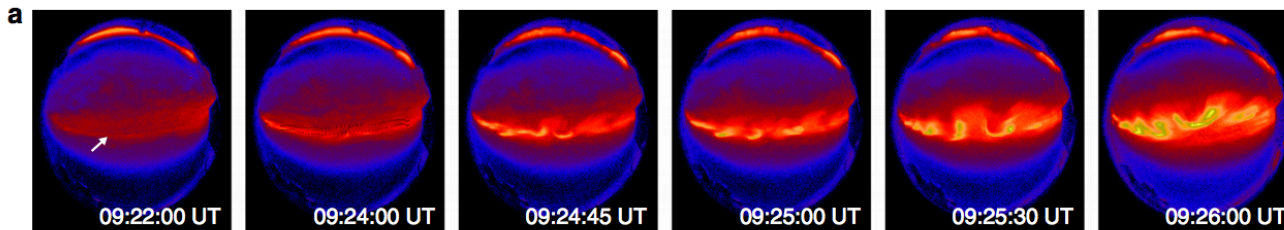
Comparison of zonal fields spatial scale



Conclusions

- Global gyrokinetic simulation results of nonlinear KBM show that KBM can be nonlinearly regulated and saturated by the zonal fields, including the zonal flow and the zonal current.
- An intermediate regime resembling that discovered in full MHD simulations is observed. Current sheet thinning can destabilize tearing modes.
- At the narrow pedestal region, zonal fields shearing scale is small, and its effects can be suppressed to some extent by the non-zonal nonlinear convection.

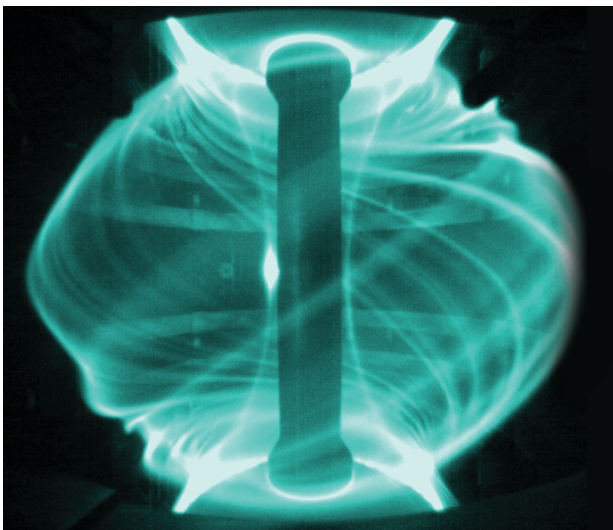
Disruptive problems are complex in nature and involve coupling of multi physical mechanisms
— — Machine Learning?



Substorm onset on Sep 2012 [Kalmoni 2018]

Avalanche?
Reconnection?
KBM?
KAW?

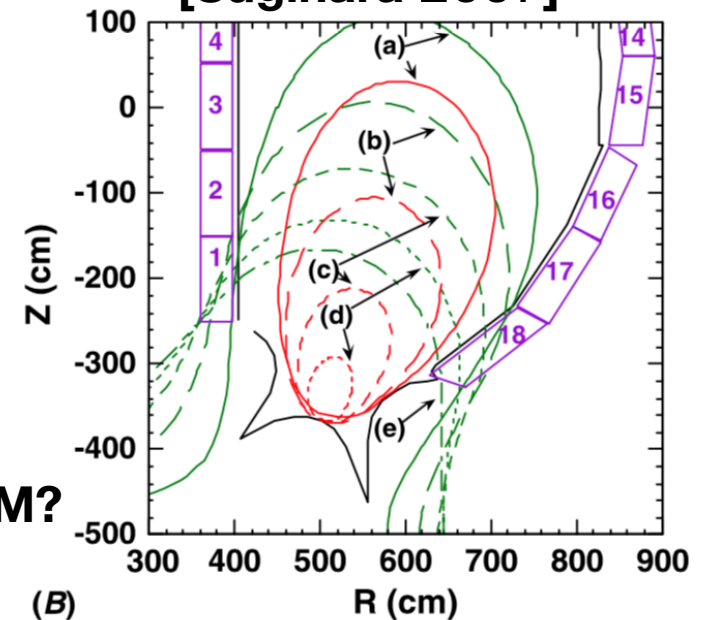
ELM Filaments in MAST



KBM?
MTM?
Resonance?

Avalanche?
NTM?
What drives NTM?

LCFS in Disruption [Sugihara 2007]





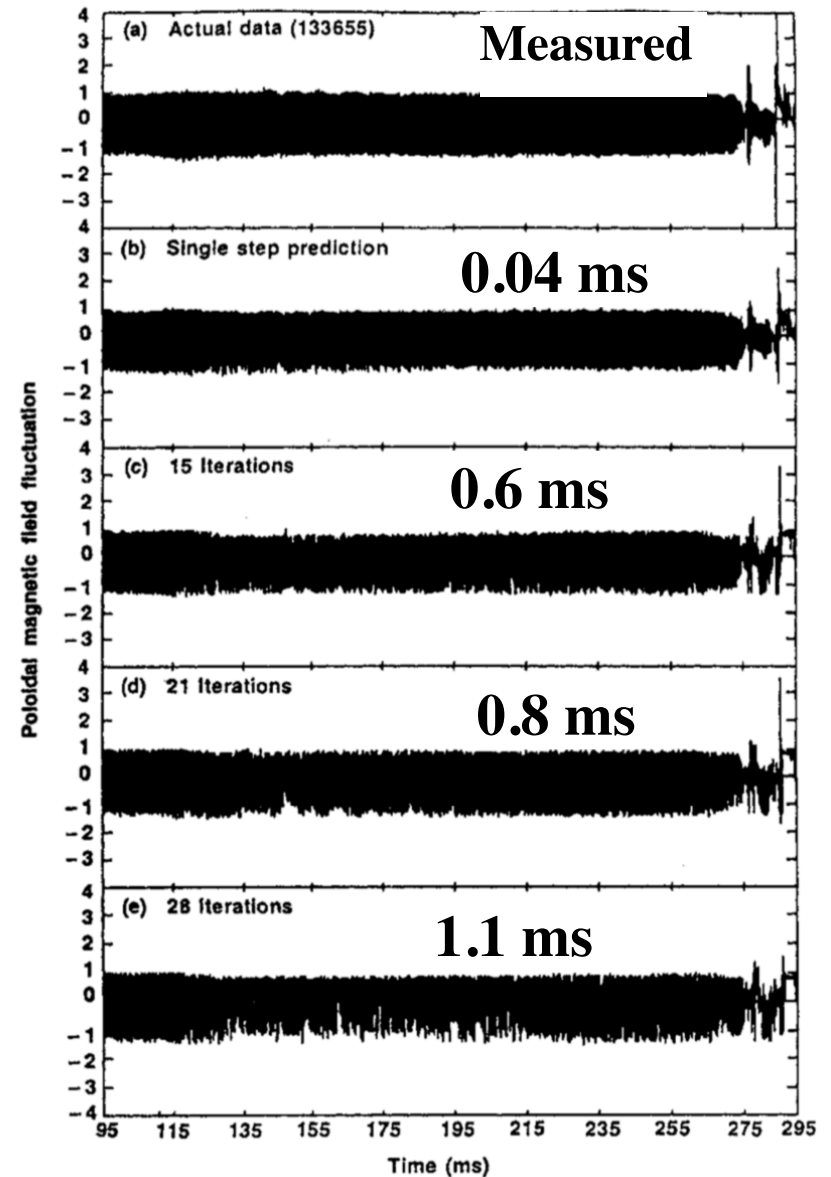
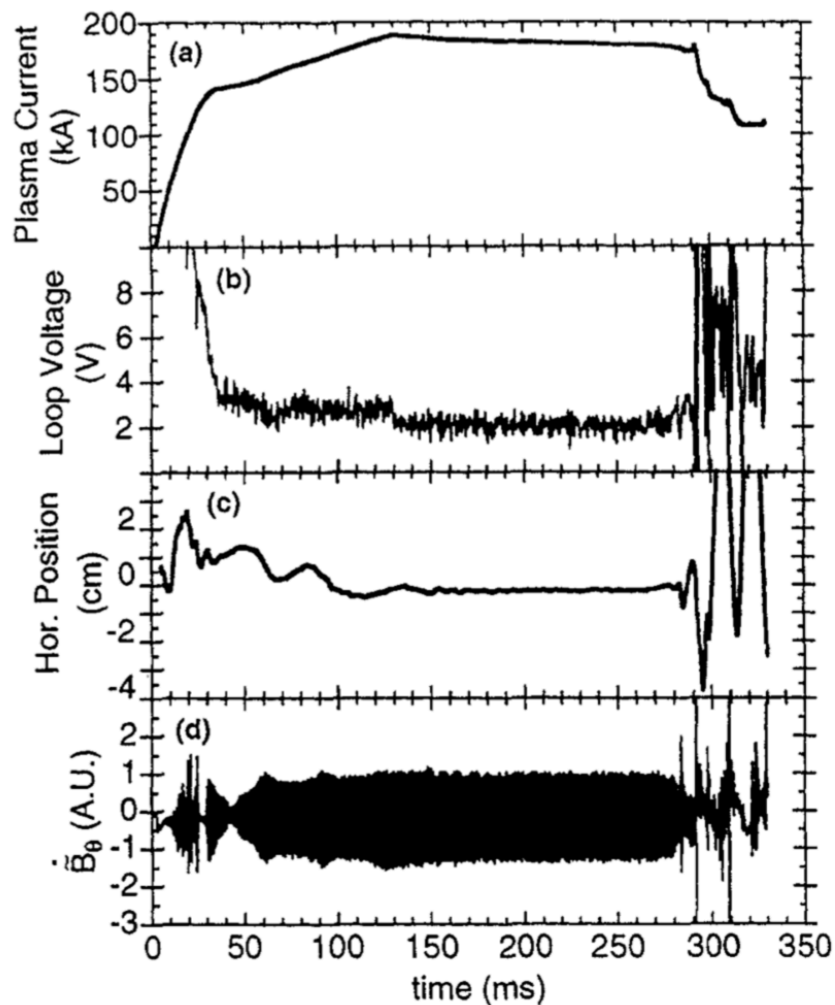
Machine Learning in disruption predictions (1996)

Effort since 90s. [J.V. Hernandez 1996]

1-2 layers of feedforward neural nets

2018: ~30 warning time, >90% Accuracy

SVM+Deep learning [Ferreira 2018]

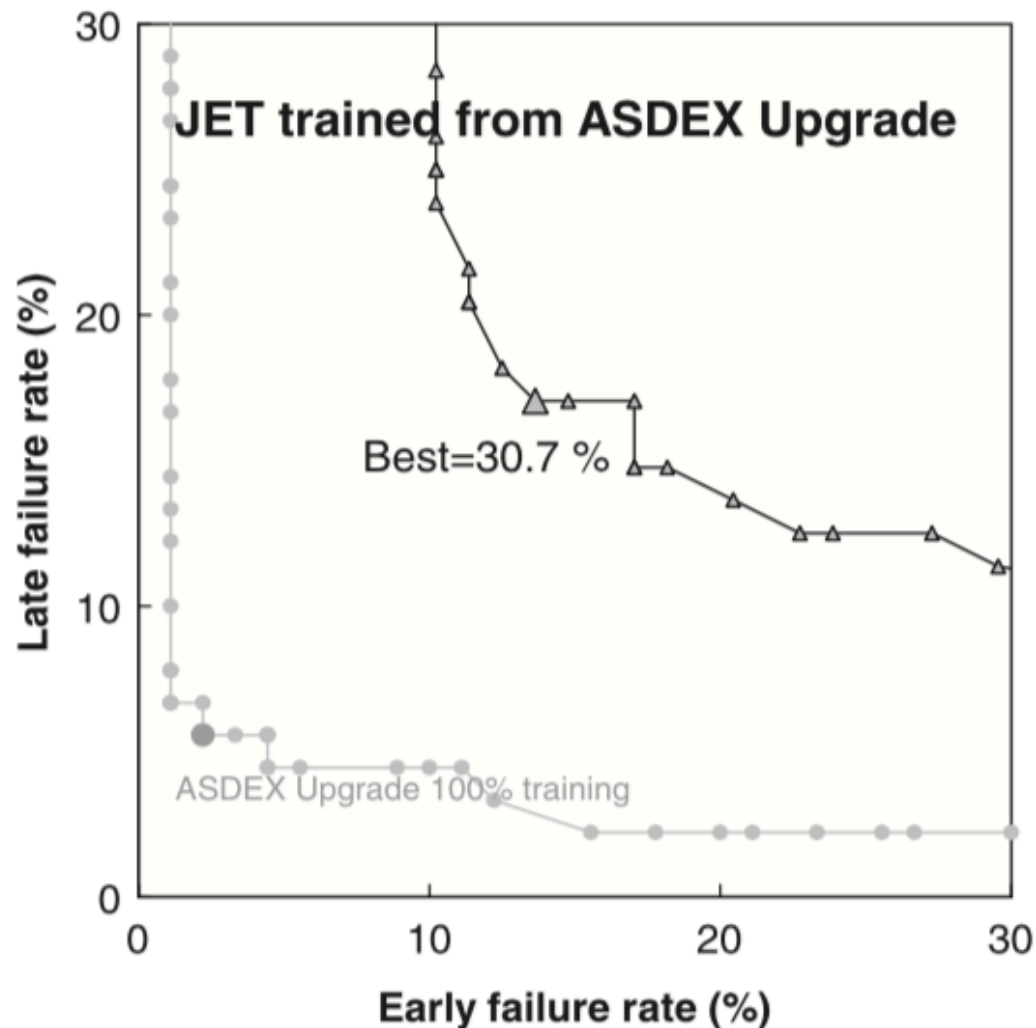


Machine Learning in disruption predictions (2005)

Cross-tokamak prediction is hard [C. Windsor 2005]

Two layer Neural Network

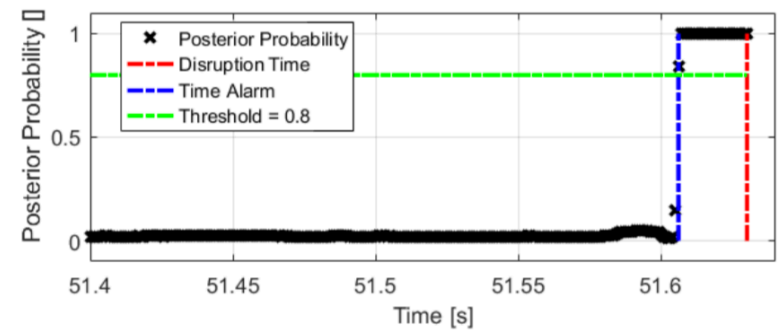
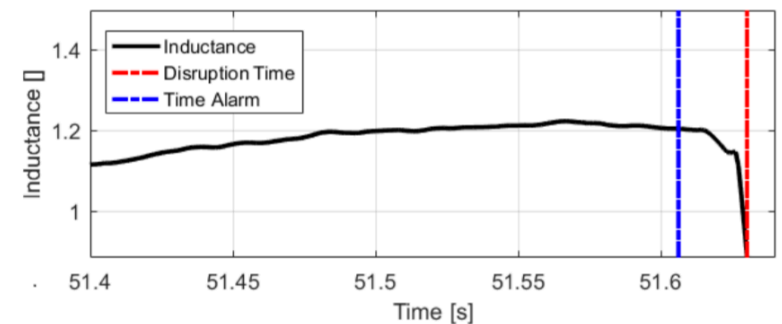
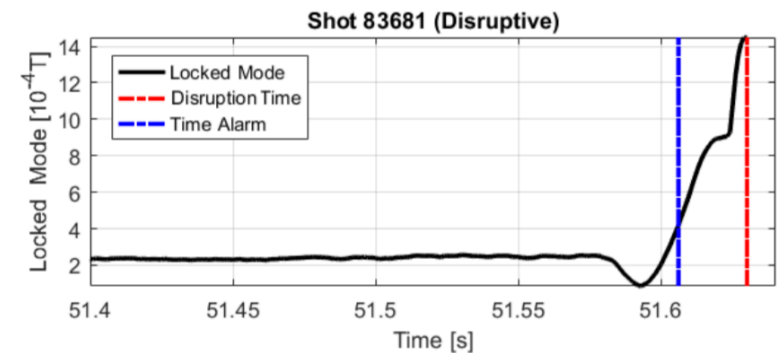
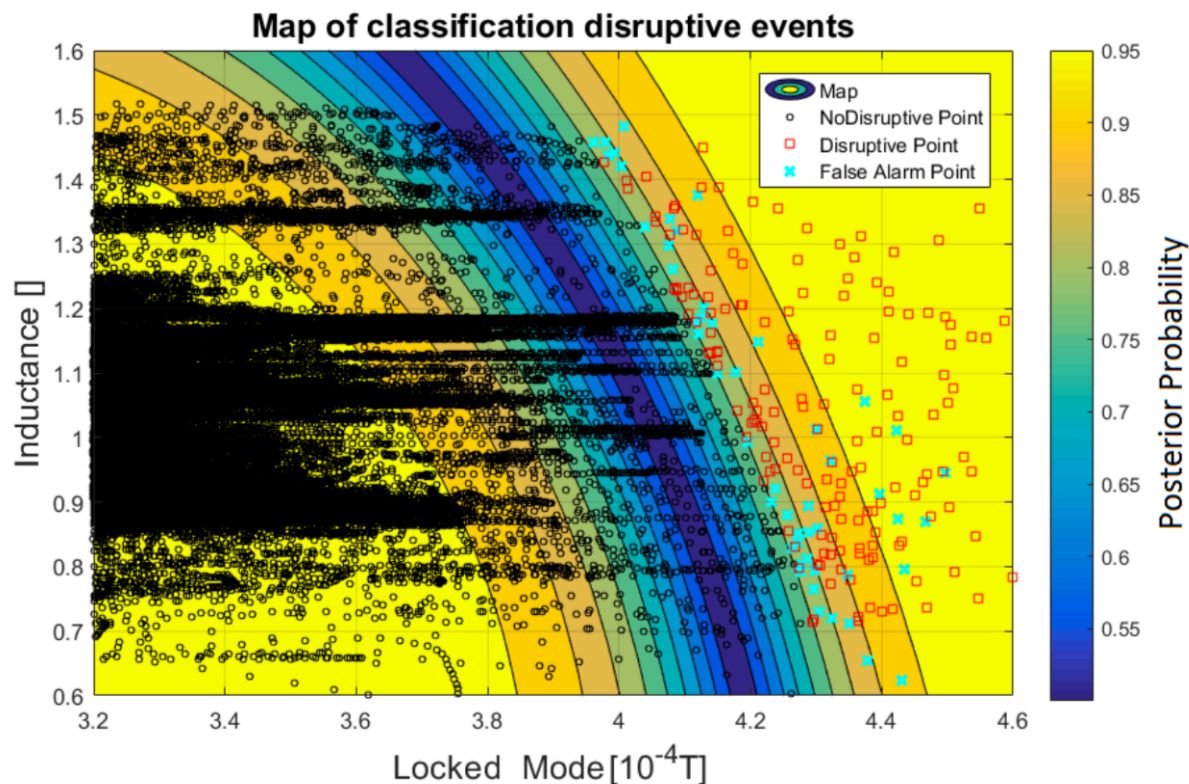
JET \longleftrightarrow ASDEX Upgrade (< 70%, 10ms)



Machine Learning in disruption predictions (2018)

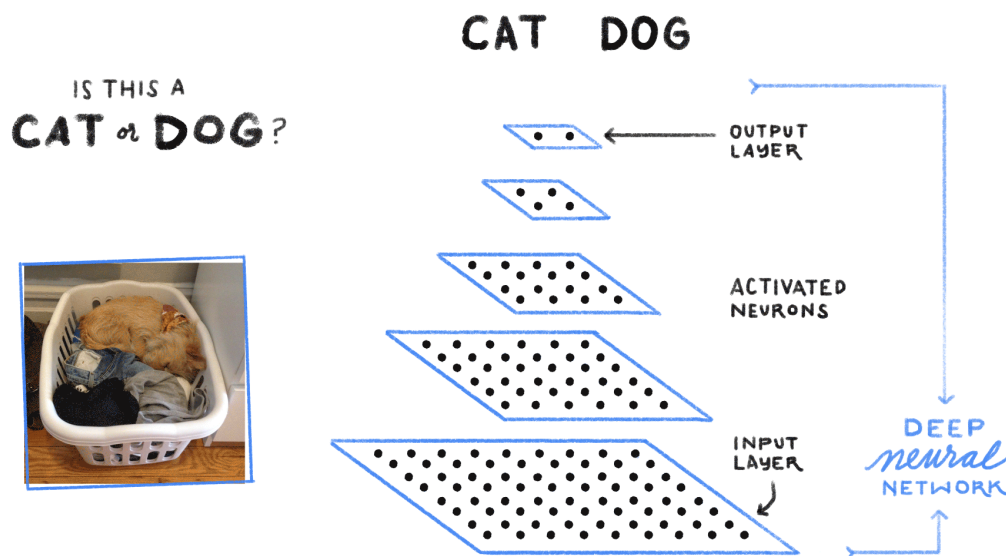
Some 'adaptivity' achieved. [A. Murari 2018]
Probabilistic SVM

**Average warning time longer than 300ms
(~95% True Positive, 5% False Positive)**





Deep Learning vs “Shallow” Learning

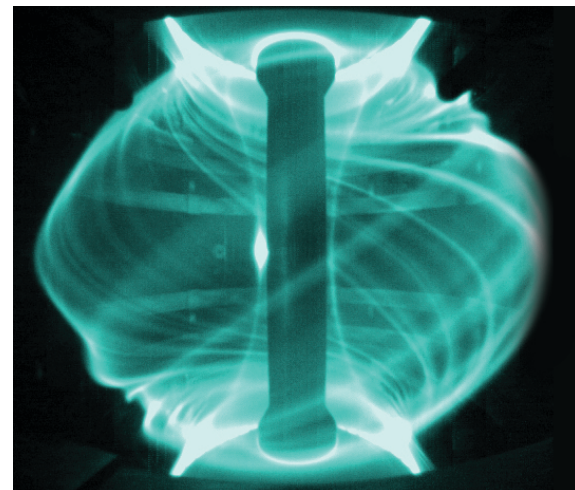
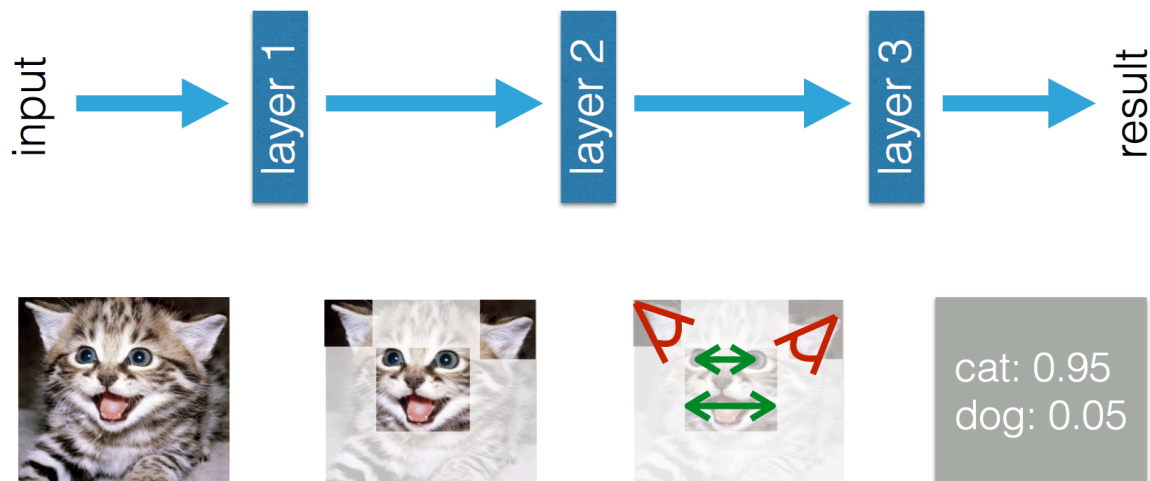


**Hierarchical representation
of Complex data**

Deep RNNs -> handwriting/ speech
recognition

Deep CNNs-> image recognition/
AlphaGo

**Q: Is this going to
disrupt?**



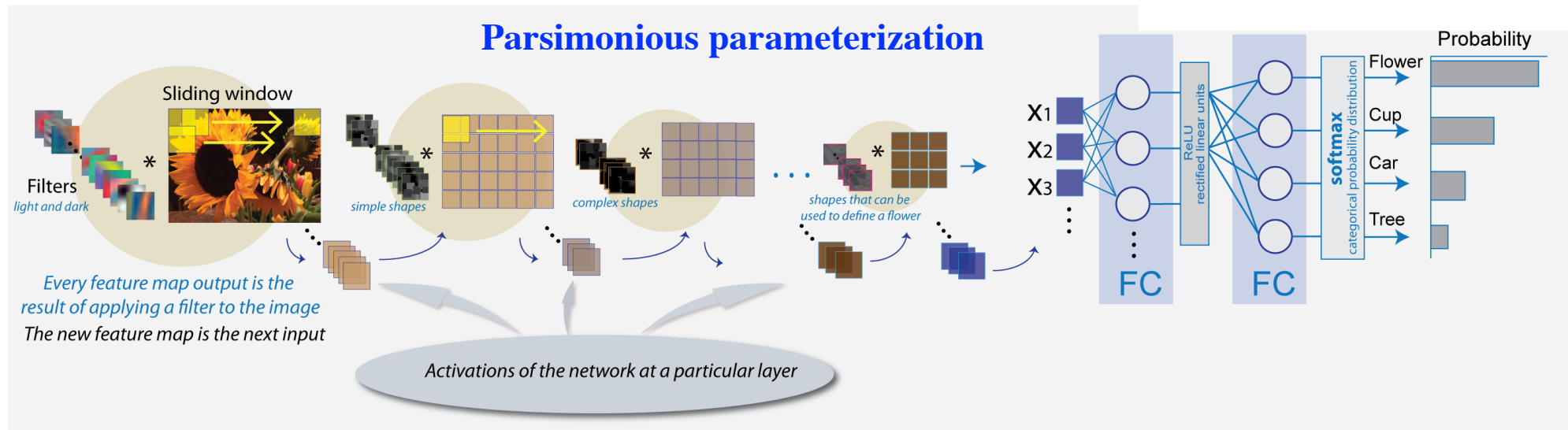


Deep Learning — CNNs

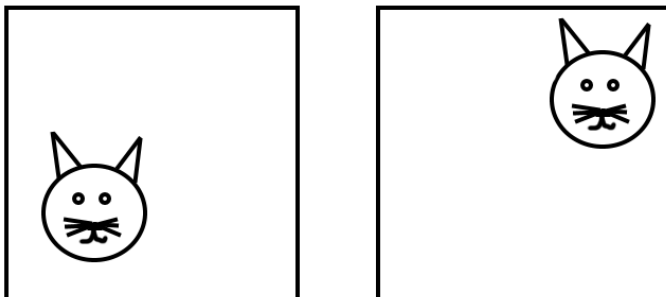
Superhuman performance of Go game
Image recognition



Example: Cifar10 data
10 classes, 60000
32X32 images
Error rate <2% (2018)

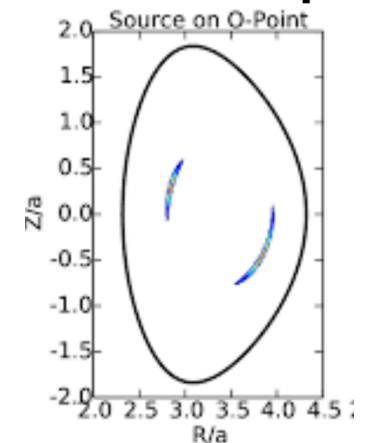
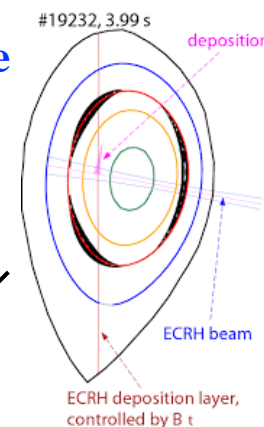


Cat is always a cat



Translational Invariance

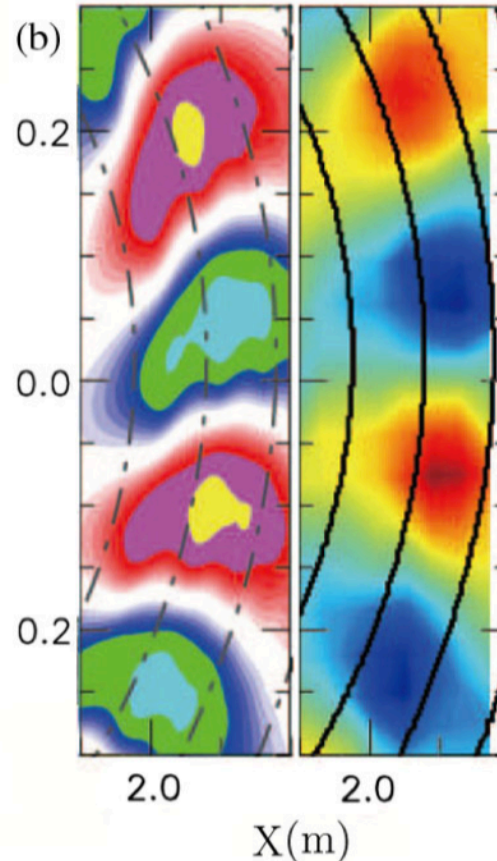
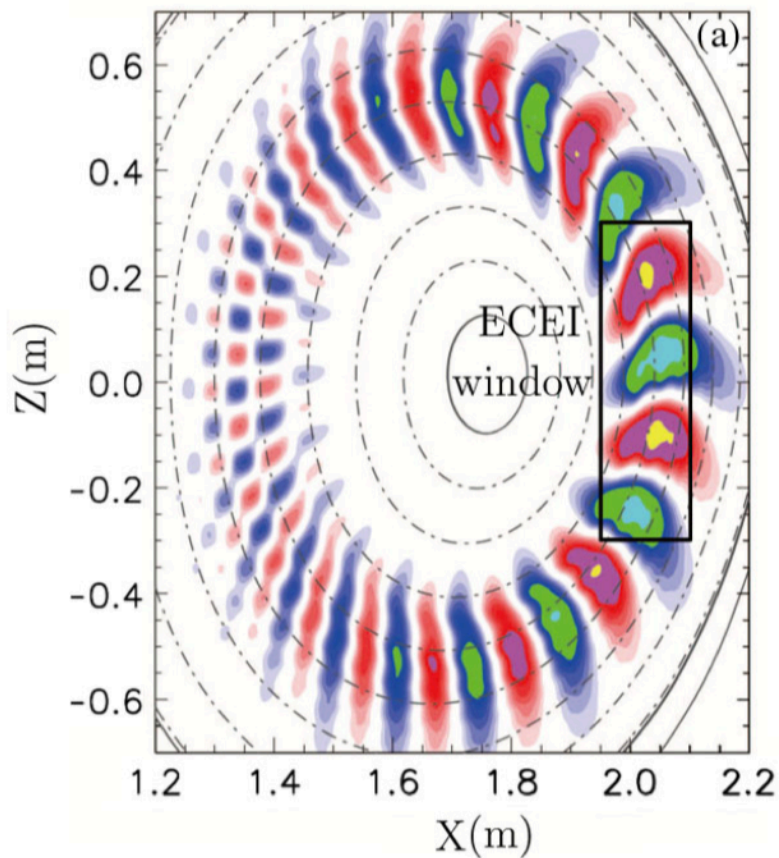
Certain island always cause disruption



Q: But how can we get these detailed data?

Q: But how can we get these detailed data?

— First-Principles-Simulation results as virtual experimental results



Linear Properties
(Useful even in simple
regression models)

Complete particle/field
information

Reproduce any
diagnostic in real/
velocity space

Perfect 2/3 dimensional
data for Neural
Networks

Comparison of TAE Te structure from the simulation (left) and from the DIII-D experiment (right) in the ECEI window [Z. Wang, 2013]

1. *Journal of the American Medical Association*, 2000; 284: 2689-2695.

What factor, which physics is more important to disruption?



Fix imbalanced data set

**Select sensitive parameters/
Extract useful phase space info
Reduce particle Noise**

```

! physical unit for equilibrium
! 1: Input reference flux surface (iflux) values for etemp0 and eden0
etemp0=196.9 !150 !287.29 ! on-axis electron temperature, unit=ev
eden0=1.496e14 !0.21e14 !0.4178e14 ! on-axis electron number density, unit=1/cm^3
r0=177.1 ! major radius, unit=cm
b0=16793.0 ! on-axis magnetic field, unit=gauss
/

%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%

&Equilibrium_parameters

! example numerical/analytical equilibrium using variables from numreq=1

psiw_analytic= 3.75e-2 ! poloidal flux at wall
ped_analytic= 3.75e-2 ! poloidal flux at separatrix

! q and zeff profile is parabolic: q=q1+q2*psi/(psiw+q3)*(psi/psiw)^2
q_analytic= 0.8200 1.10000 1.00000 !
ze_analytic= 1.0 0.0 0.0 !
er_analytic= -0.06312057 0.1661067 0.0 !

itemp0_analytic= 1.0 ! on-axis thermal ion temperature, unit=T_e0
fitemp0_analytic= 2.0 ! on-axis fast ion temperature, unit=T_e0
fnden0_analytic= 1.0e-5 ! on-axis fast ion density, unit=n_e0
fitemp0_analytic= 1.0 ! on-axis fast electron temperature, unit=T_e0
feden0_analytic= 1.0 ! on-axis fast electron density, unit=n_e0

```




Summary

- Deep learning could achieve break through in disruption research, with the help from first-principles based codes.
- The linear properties as well as nonlinear mode structures from the simulations could be incorporated into the deep learning models for disruption predictions in the form of a new parameter/channel, as a first-principles physics guide to the AI.
- The deep learning model could in turn provide feedback on the sensitivity of the parameters and thus automatically select new inputs for the first-principles codes.

Thank you!

**Following Slides are
Supplementary materials**

Gyrokinetic simulation model

GTC Flow Chart-Conservative

$$\begin{aligned} \frac{dw_\alpha}{dt} = (1 - w_\alpha) & \left[- \left(v_\parallel \frac{\delta \mathbf{B}_\perp}{B_\parallel^*} + \mathbf{v}_E + \mathbf{v}_{b\parallel} + \mathbf{v}_{NL} \right) \cdot \boldsymbol{\kappa}_\alpha + \frac{u_{\parallel\alpha 0} \delta \mathbf{B}_\perp}{T_{\alpha 0} B_0} \cdot \nabla (\mu B_0 + Z_\alpha \Phi_{eq}) \right. \\ & - \frac{v_\parallel \delta \mathbf{B}_\perp}{T_{\alpha 0} B_0} \cdot \nabla (Z_\alpha \Phi_{eq}) - \left[\mathbf{b}_0 \cdot \nabla (Z_\alpha \phi + Z_\alpha \phi_{NL} + \mu \delta B_\parallel) + \frac{\partial \delta A_\parallel}{\partial t} \right] \frac{Z_\alpha (v_\parallel - u_{\parallel\alpha 0})}{T_{\alpha 0}} \\ & \left. - \frac{1}{T_{0\alpha}} \left(\mathbf{v}_g + \left(\frac{v_\parallel \delta \mathbf{B}_\perp}{B_\parallel^*} + \mathbf{v}_c + \mathbf{v}_{deq} \right) (1 - \frac{u_{\parallel\alpha 0}}{v_\parallel}) \right) \cdot \nabla (Z_\alpha \phi + Z_\alpha \phi_{NL} + \mu \delta B_\parallel) \right] \end{aligned}$$

$$\begin{aligned} \frac{\partial \delta n_e}{\partial t} + \mathbf{B}_0 \cdot \nabla \left(\frac{n_0 \delta u_{\parallel e}}{B_0} \right) + B_0 \mathbf{v}_E \cdot \nabla \left(\frac{n_0}{B_0} \right) - n_0 (\mathbf{v}_* + \mathbf{v}_E) \cdot \frac{\nabla B_0}{B_0} \\ + \delta \mathbf{B}_\perp \cdot \nabla \left(\frac{n_0 u_{\parallel 0}}{B_0} \right) + \frac{c \nabla \times \mathbf{B}_0}{e B_0^2} \left(-\nabla \delta P_\parallel - \frac{(\delta P_\perp - \delta P_\parallel) \nabla B_0}{B_0} + n_0 e \nabla \delta \phi \right) \\ + \delta \mathbf{B}_\perp \cdot \nabla \left(\frac{n_0 e \delta u_{\parallel e}}{B_0} \right) + B_0 \mathbf{v}_E \cdot \nabla \left(\frac{\delta n_e}{B_0} \right) + \frac{c \delta n_e}{B_0^2} \mathbf{b}_0 \times \nabla B_0 \cdot \nabla \phi + \frac{c \delta n_e}{B_0^2} \nabla \times B_0 \cdot \nabla \phi \\ + \frac{c \mathbf{b}_0 \times \nabla \delta B_\parallel}{e} \nabla \left(\frac{\delta P_\perp + P_{\perp 0}}{B_0^2} \right) + \frac{c \nabla \times \mathbf{b}_0 \cdot \nabla \delta B_\parallel}{e B_0^2} (\delta P_\perp + P_{\perp 0}) = 0 \end{aligned}$$

$$\frac{\partial \delta A_\parallel^A}{\partial t} = c \mathbf{b}_0 \cdot \nabla \phi_{ind}, \quad \delta n_e^{(a)} = n_0 \left(\frac{e \delta \phi + e \delta \phi_{ind}}{T_e} - \frac{\delta B_\parallel}{B_0} \right) + \frac{\partial n_0}{\partial \psi_0} \delta \psi + \frac{\partial n_0}{\partial \alpha_0} \delta \alpha$$

$$\left(\nabla_\perp^2 - \frac{\omega_{pe}^2}{c^2} \right) \frac{\partial \delta A_\parallel^{NA}}{\partial t} = \frac{\omega_{pe}^2}{c} \chi_\parallel - c \nabla_\perp^2 (\mathbf{b}_0 \cdot \nabla \phi_{ind})$$

$$\frac{e \delta \phi}{T} = (1 - \rho_s^{-2} \nabla_\perp^2) \frac{\delta n_i - \delta n_e}{n_0}$$

$$\frac{\delta B_\parallel}{B_0} = \frac{\beta}{2} \left[-\frac{\delta P_e + \delta \tilde{P}_i}{P_0} + \frac{\delta n_i - \delta n_e}{n_0} + (1 - \rho_s^2 \nabla_\perp^2)^{-1} \frac{\delta n_i - \delta n_e}{n_0} \right]$$

$$\delta f_i^n, \delta n_e^n, \delta \phi^n, \delta \phi_{ind}^n, \delta A_\parallel^{A,n}, \delta A_\parallel^{NA,n}, \delta B_\parallel^n, \delta h_e^n, \delta u_{e\parallel}^n$$

$$\begin{aligned} \delta f_i^{n+1} &= \frac{\partial f_i}{\partial t} (\delta \phi^n, \delta \phi_{ind}^n, \delta B_\parallel^n, \delta A_\parallel^{A,n}, \delta A_\parallel^{NA,n}) \Delta t + \delta f_i^n \\ \delta n_e^{n+1} &= \frac{\partial \delta n_e}{\partial t} (\delta u_{e\parallel}^n, \delta \phi^n, \delta \phi_{ind}^n, \delta B_\parallel^n, \delta A_\parallel^{A,n}, \delta A_\parallel^{NA,n}) \Delta t + \delta n_e^n \\ \delta A_\parallel^{A,n+1} &= \frac{\partial \delta A_\parallel^A}{\partial t} (\delta \phi_{ind}^n) \Delta t + \delta A_\parallel^{A,n} \\ \delta A_\parallel^{NA,n+1} &= \frac{\partial \delta A_\parallel^{NA}}{\partial t} (\delta u_{e\parallel}^n, \delta \phi^n, \delta \phi_{ind}^n, \delta B_\parallel^n, \delta A_\parallel^{A,n}, \delta A_\parallel^{NA,n}) \Delta t + \delta A_\parallel^{NA,n} \\ \delta h_e^{n+1} &= \frac{\partial \delta h_e}{\partial t} \left(\frac{\partial \delta f_e^{(a),n}}{\partial t}, \delta \phi^n, \delta \phi_{ind}^n, \delta B_\parallel^n, \delta A_\parallel^{A,n}, \delta A_\parallel^{NA,n} \right) \Delta t + \delta h_e^n \end{aligned}$$

$$\delta \phi_{ind}^{n+1} = \delta \phi_{ind} (\delta n_e^{n+1}, \delta \phi^{n+1}, \delta A_\parallel^{A,n+1})$$

$$\frac{\delta f^{(a),n+1}}{\partial t} = \frac{\delta f^{(a)}}{\partial t} \left(\frac{\partial \delta n_e^{n+1}}{\partial t}, \delta \phi_{ind}^{n+1}, \frac{\partial \delta B_\parallel^{n+1}}{\partial t} \right)$$

$$\delta u_{e\parallel}^{n+1} = \delta u_{e\parallel} (\delta A_\parallel^{A,n+1}, \delta A_\parallel^{NA,n+1}, \delta u_{e\parallel}^{n+1})$$

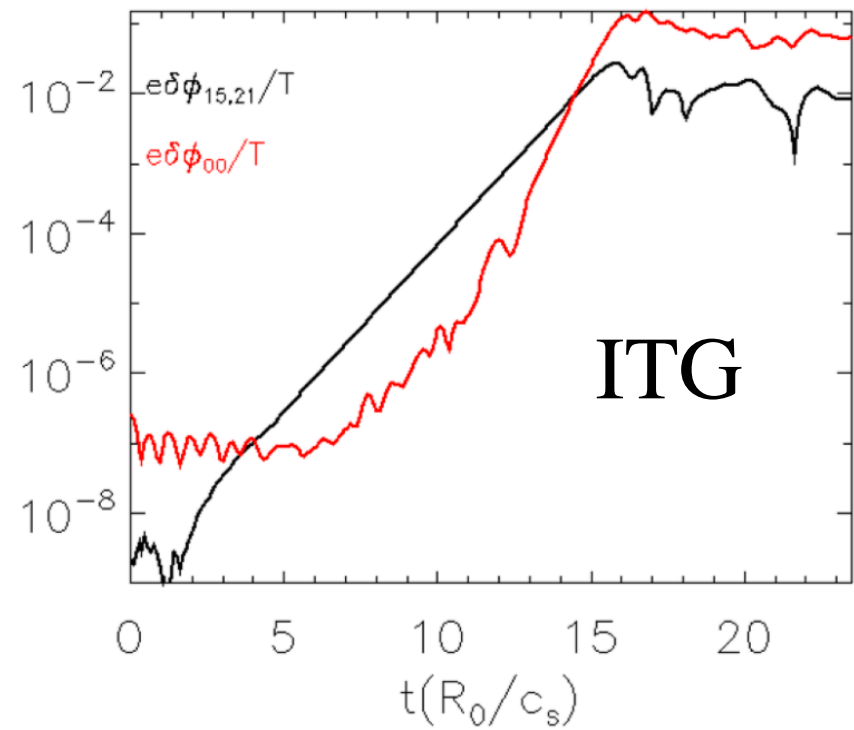
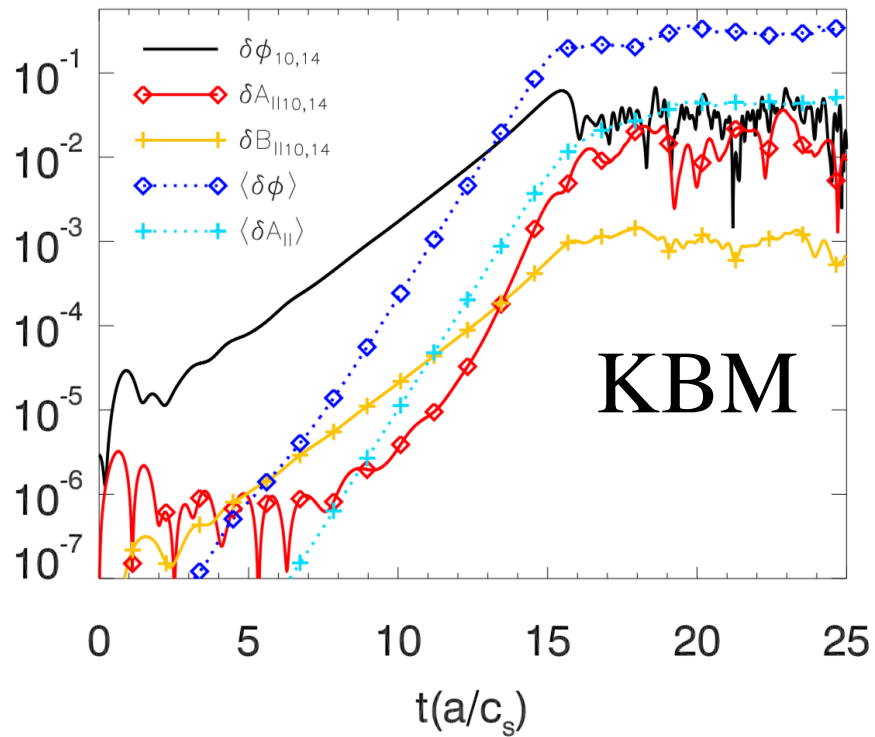
$$\delta B_\parallel^{n+1} = \delta B_\parallel (\delta P_e^{n+1}, \delta P_i^{n+1}, \delta \phi^{n+1})$$

$$\delta \phi^{n+1} = \delta \phi (\delta n_e^{n+1}, \delta n_i^{n+1}, \delta B_\parallel^{n+1})$$

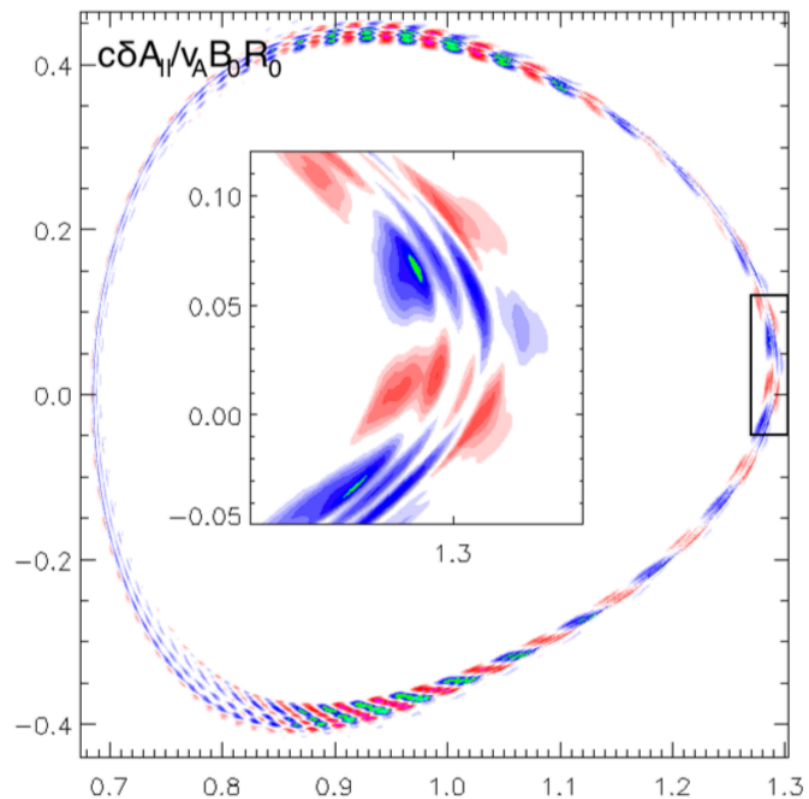
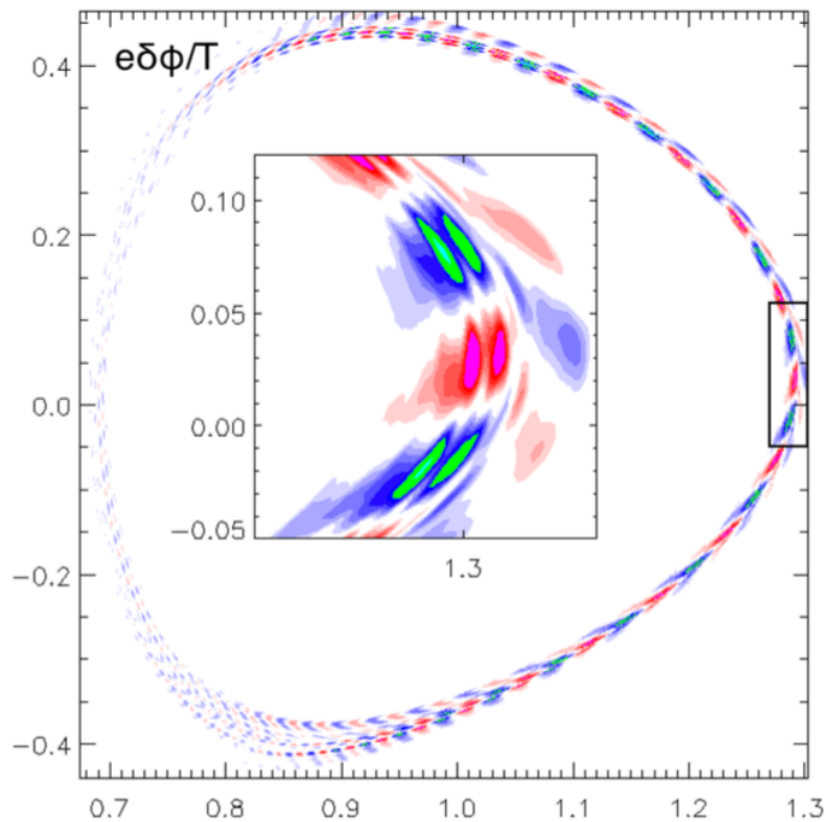
Coupled Equation

$$\delta f_i^{n+1}, \delta n_e^{n+1}, \delta \phi^{n+1}, \delta \phi_{ind}^{n+1}, \delta A_\parallel^{A,n+1}, \delta A_\parallel^{NA,n+1}, \delta B_\parallel^{n+1}, \delta h_e^{n+1}, \delta u_{e\parallel}^{n+1}$$

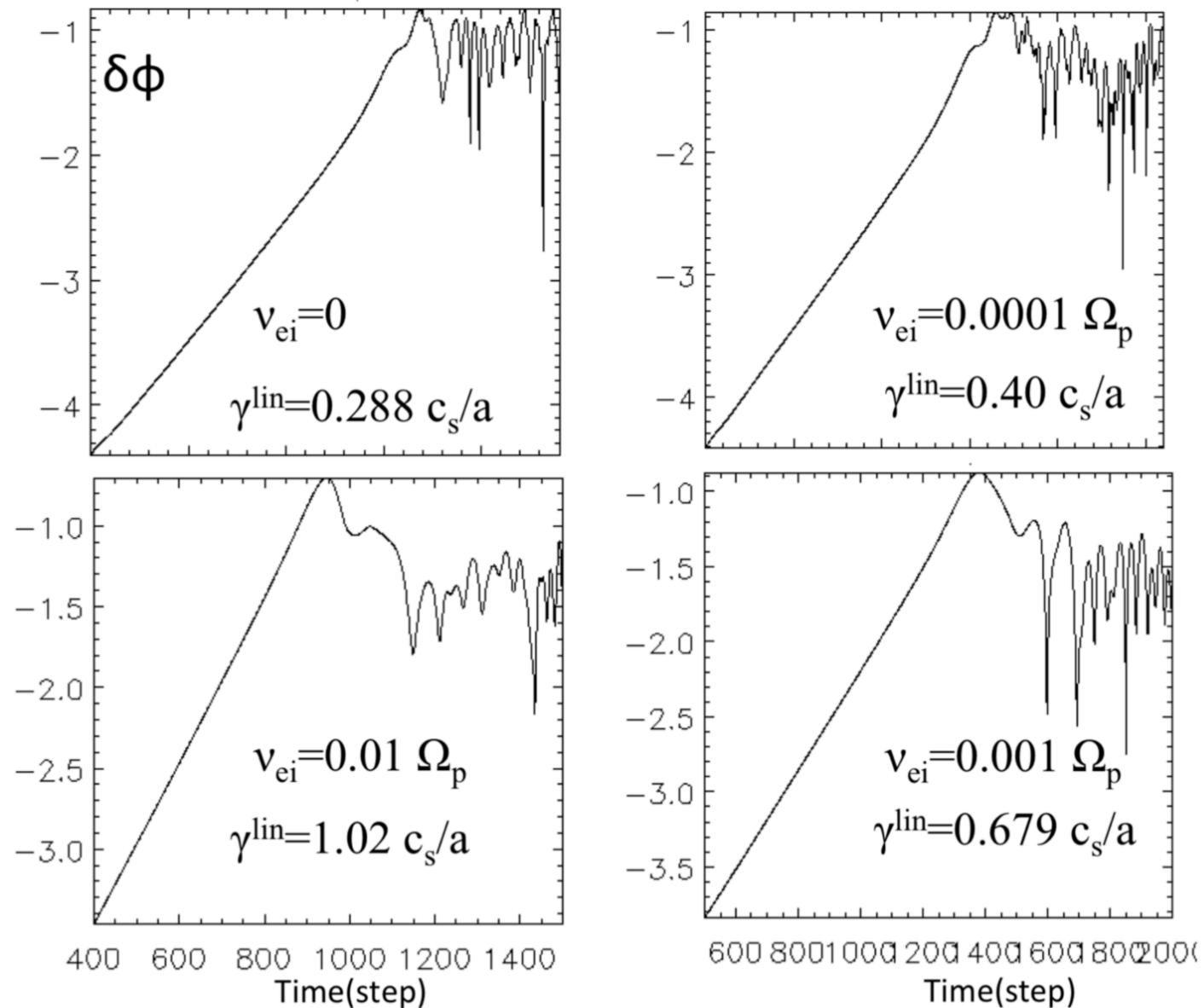
Comparison of ITG and KBM zonal flow generation



KBM nonlinear structure in later nonlinear regime in DIII-D edge



CBC case with resistivity



CBC case with resistivity

